

Magnetic Fields and Asymmetric Core Collapse Supernovae

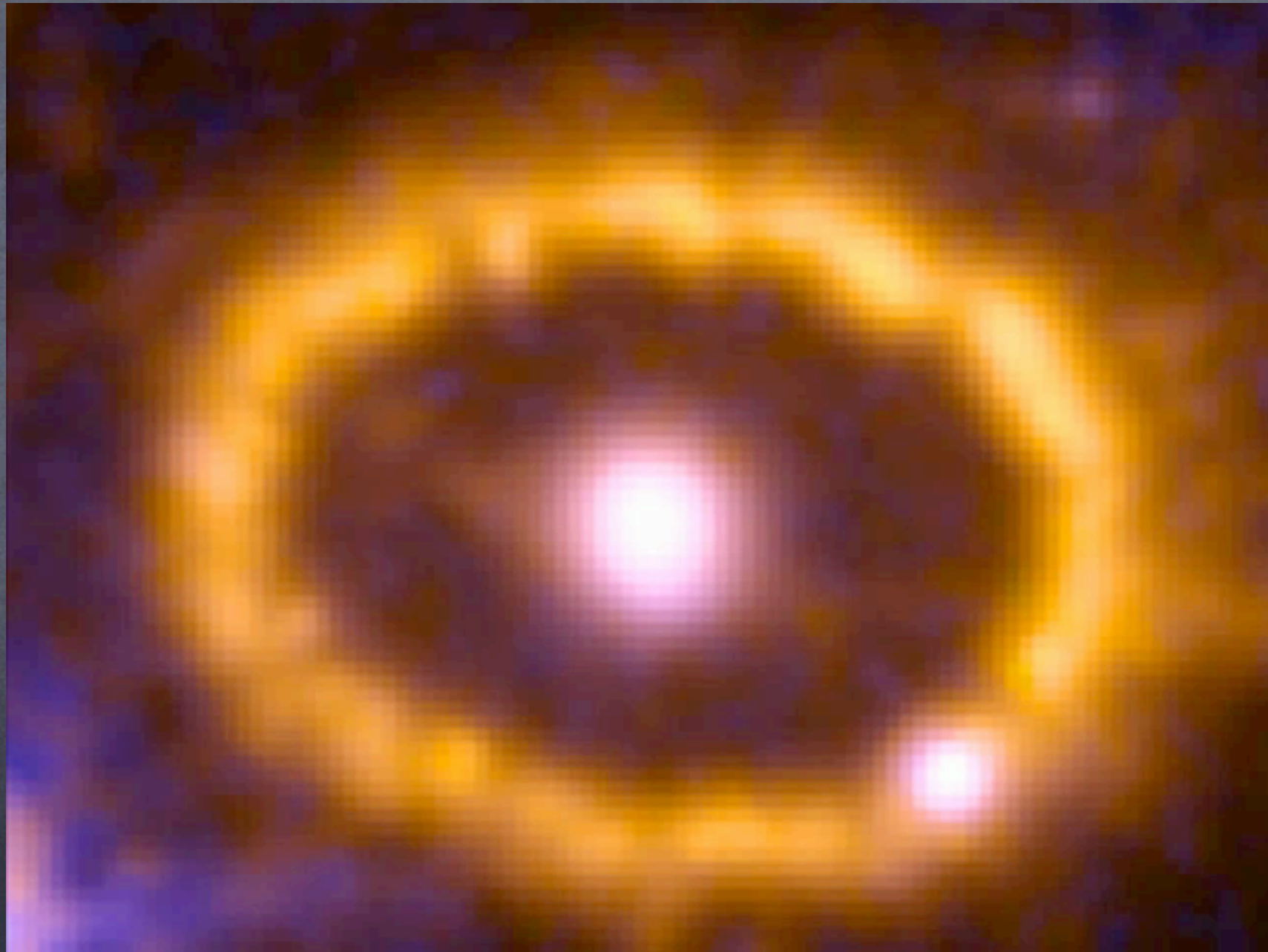
Sean M. Couch
Flash Center, University of Chicago

- CCSN Observations indicate asymmetry
- Basic Physics of CCSN & the Neutrino Mechanism
- 1D, 2D, & 3D simulations
- Magnetorotational CCSN
- Can MR-CCSN fit the observations?
- Basics of Magnetohydrodynamics
- MHD in CCSN

So, Who cares?

- Real stars rotate and are magnetic!
- The initial conditions are still uncertain...
- Could have dramatic effect on explosion dynamics
- May be critical to explaining observations
- Magnetars, Pulsars, GRBs, oh My!

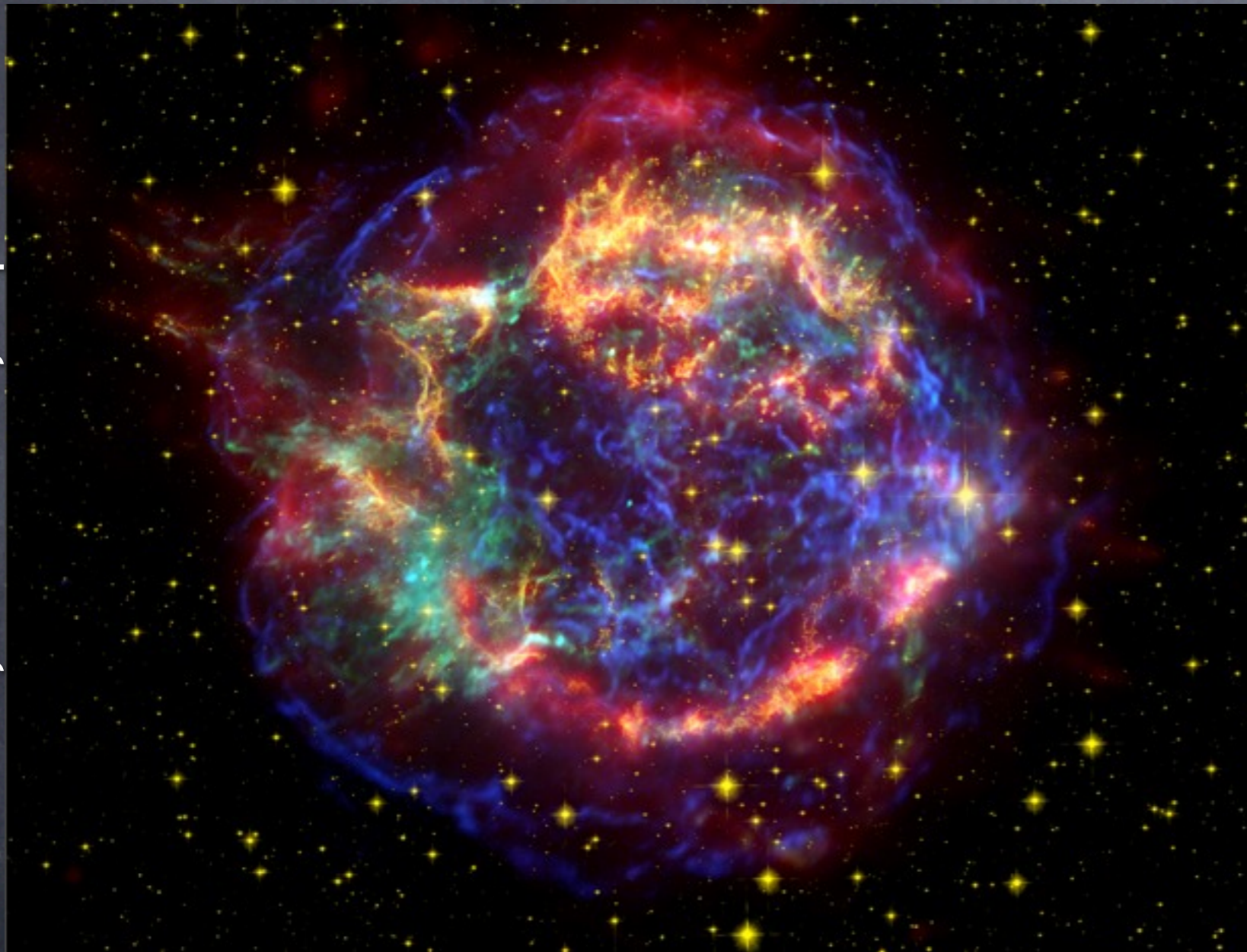
SN 1987A



Hubble

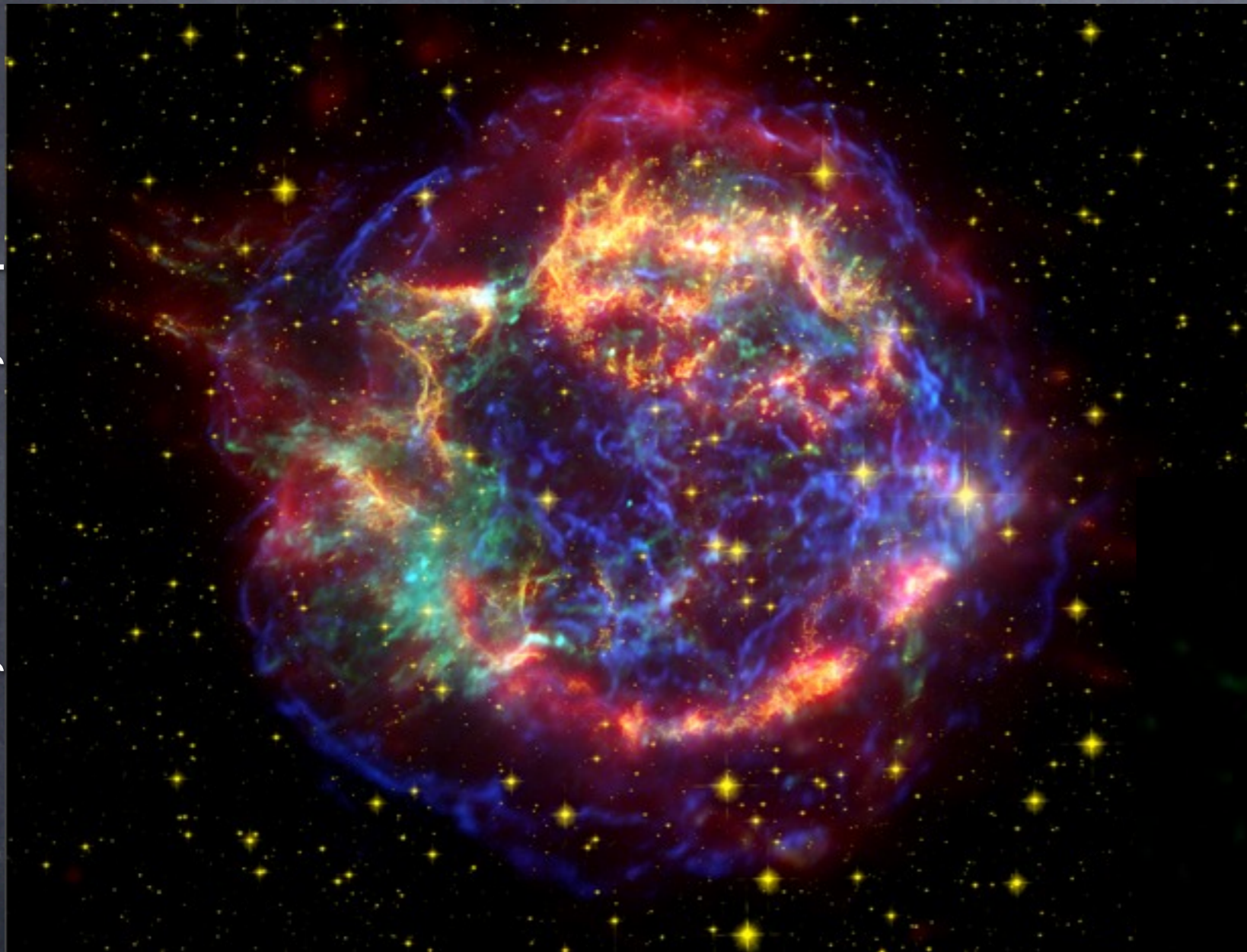
Cas A

Hubble, Chandra, Spitzer



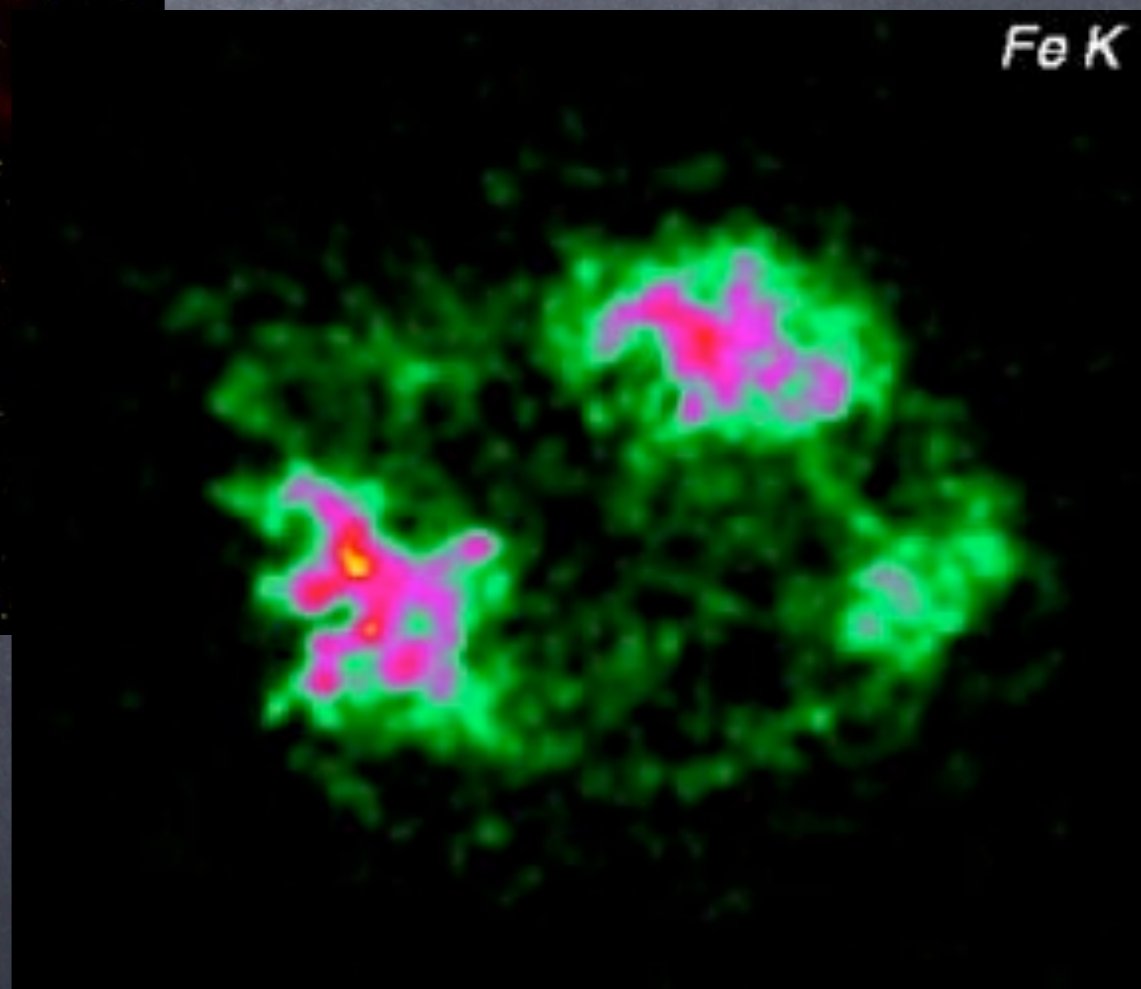
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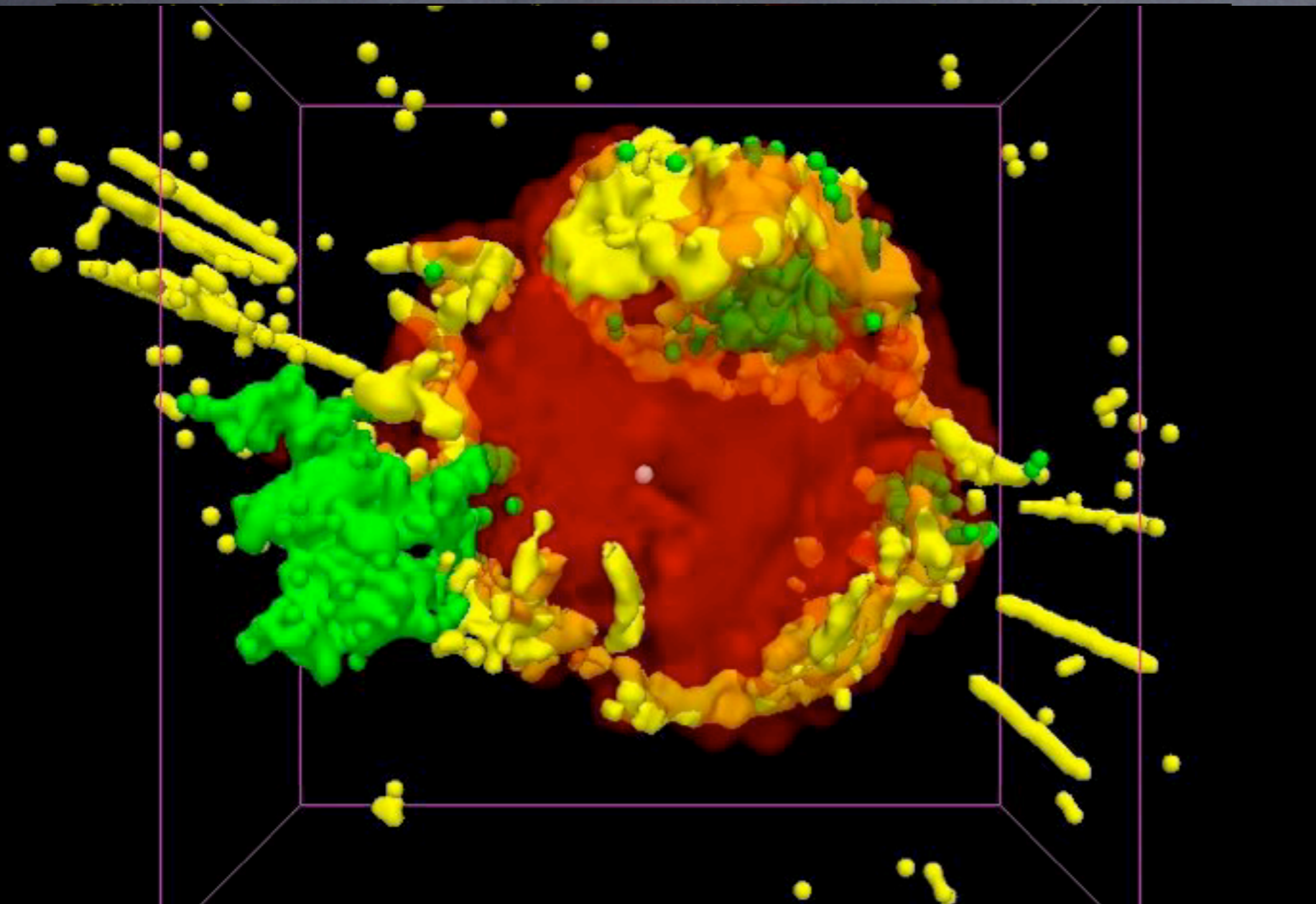


Chandra

Fe K

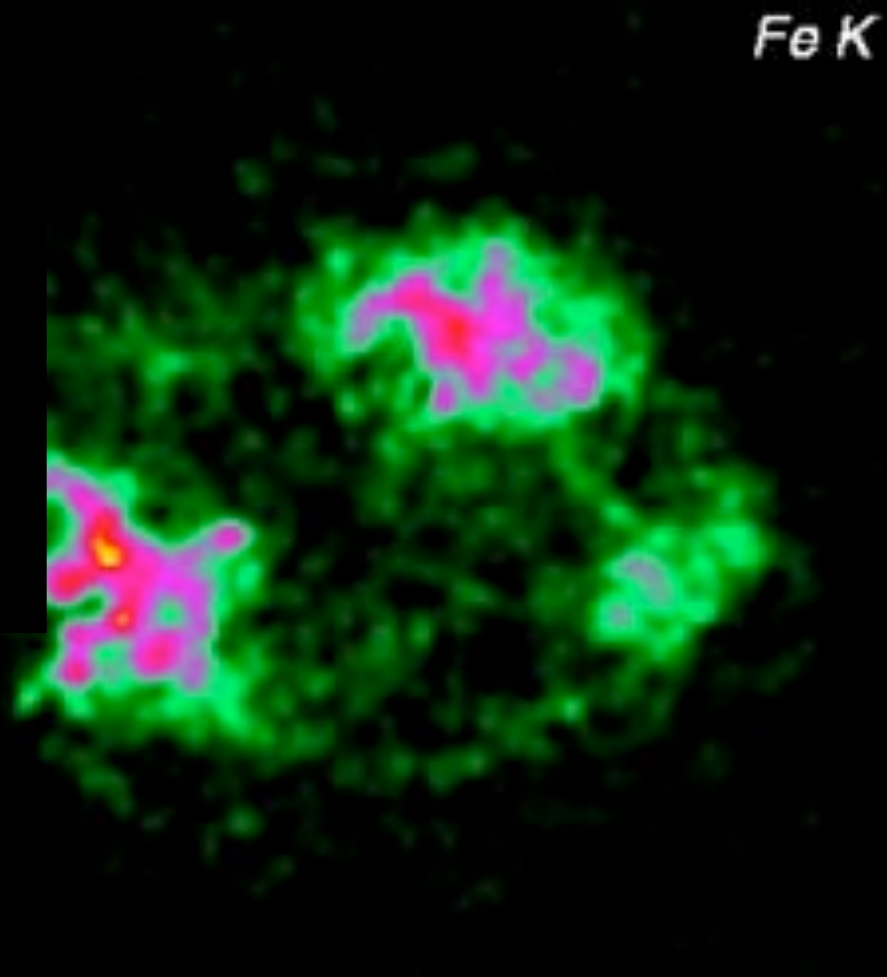


Cas A



Chandra

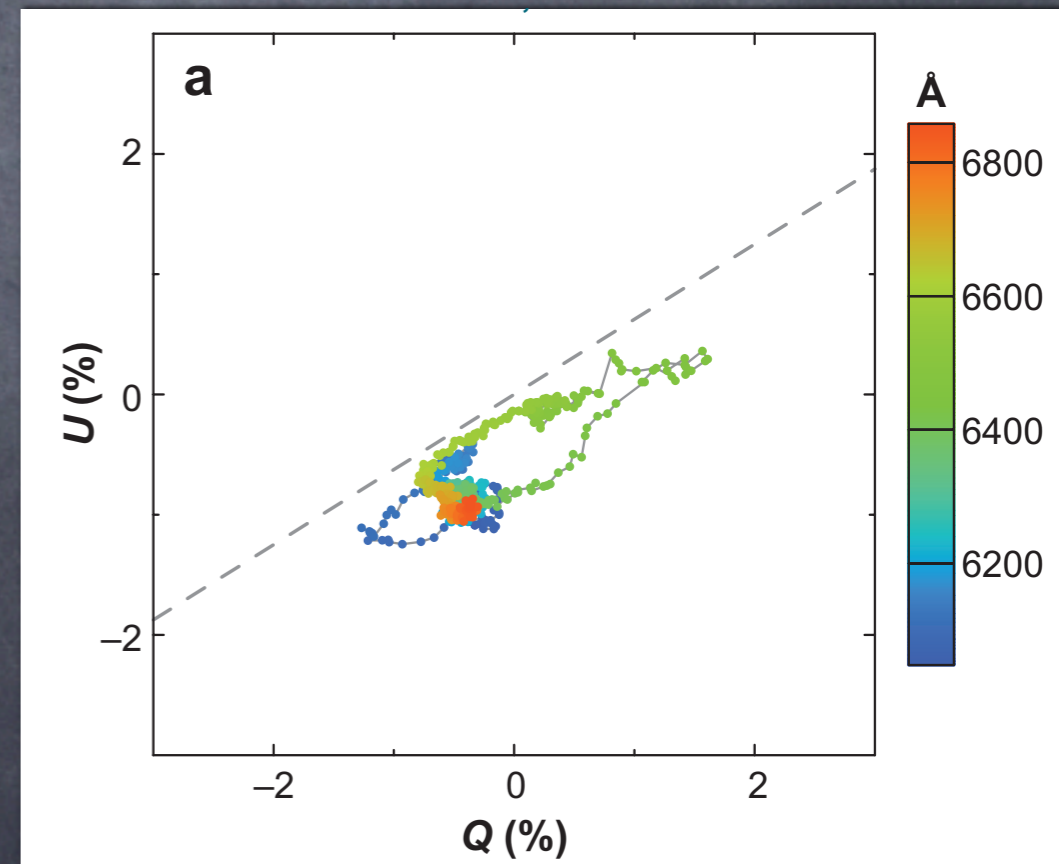
Fe K



DeLaney et al.

SN Polarization

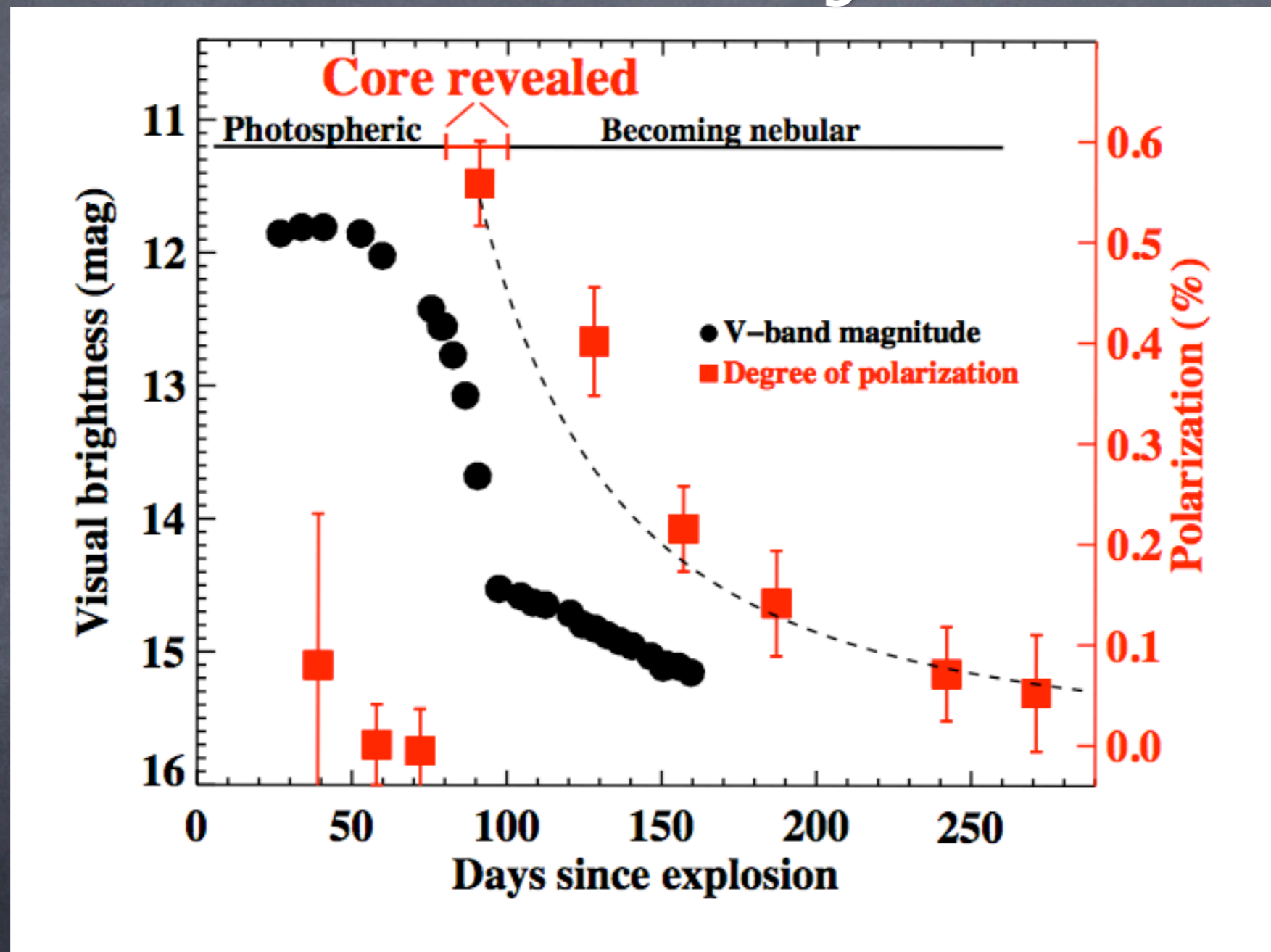
- ALL core-collapse SNe are polarized
- Higher asymmetries in the cores of explosions
- Often show a “dominant axis” in Q/U plane - indicates an elongated explosion
- Loops in Q/U plane indicate non-axisymmetry



Wang & Wheeler 2008

Type IIP Polarization

SN 2004dj

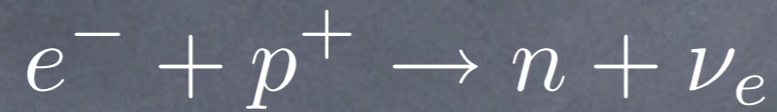


Leonard et al. 2006

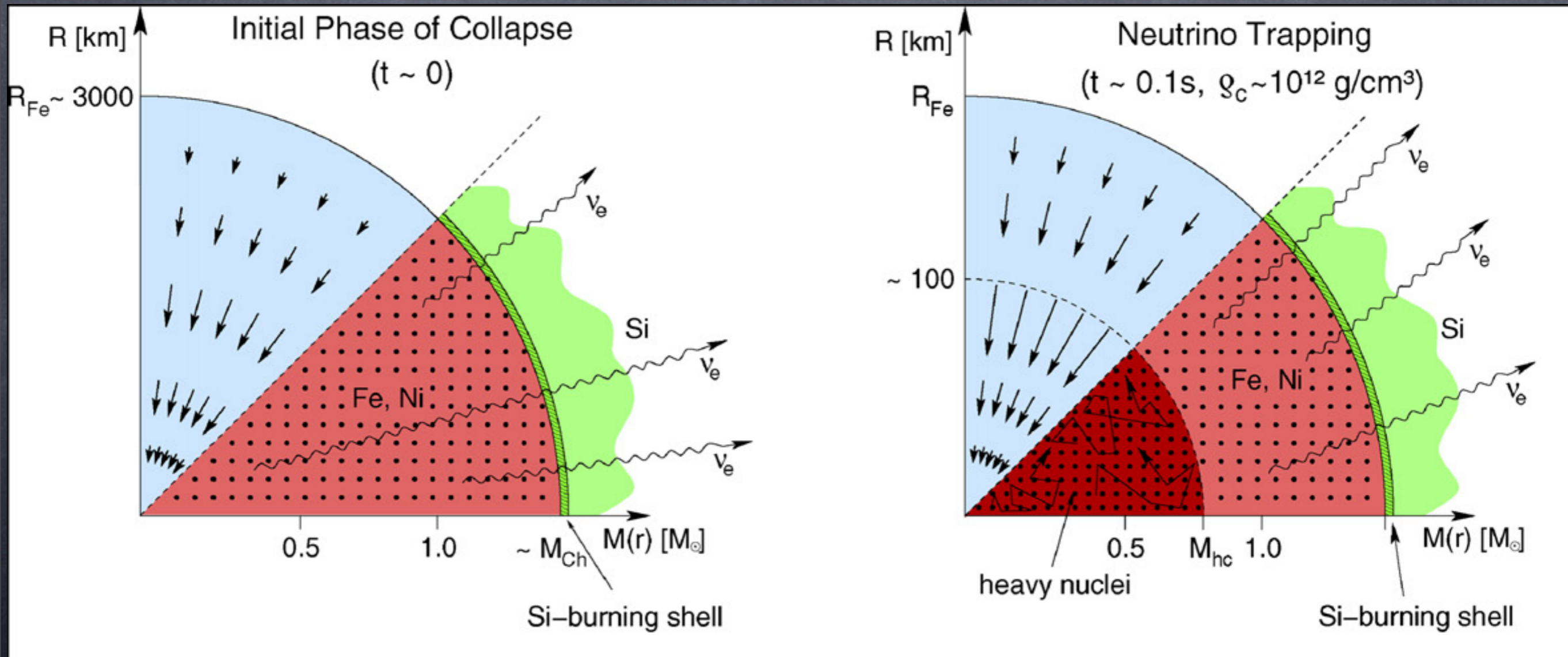
What Do the Observations Tell Us?

- Massive stars explode all the time, with energies around 10^{51} erg!
- They are NOT spherically-symmetric
- They often show general 'bi-polarity' with significant non-axisymmetry and time-dependent polarization.
- They leave remnants that often have high kick velocities and strong magnetic fields.
- Some CCSNe are associated with GRBs.
- Mixing & overturn commonly indicated.

Core Collapse Basics



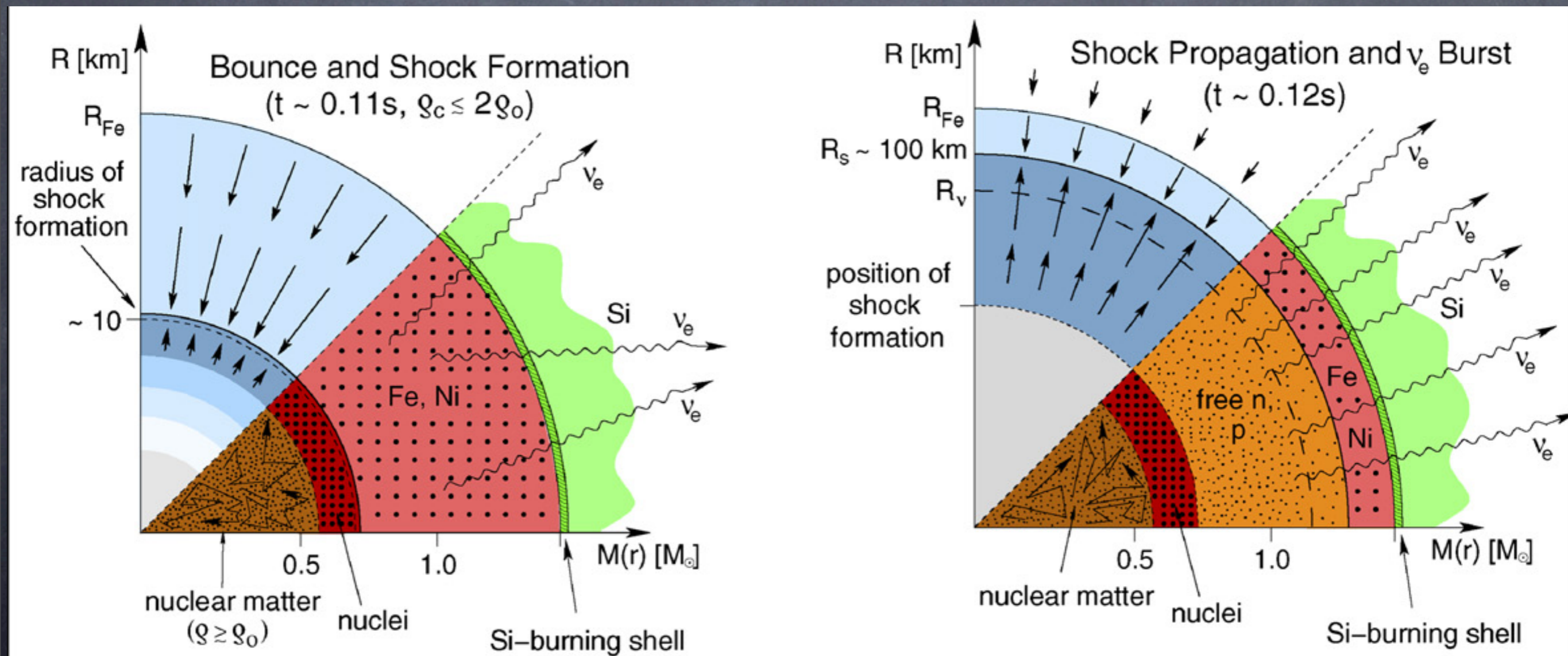
e^{-} degeneracy
pressure fails



Janka (2007)

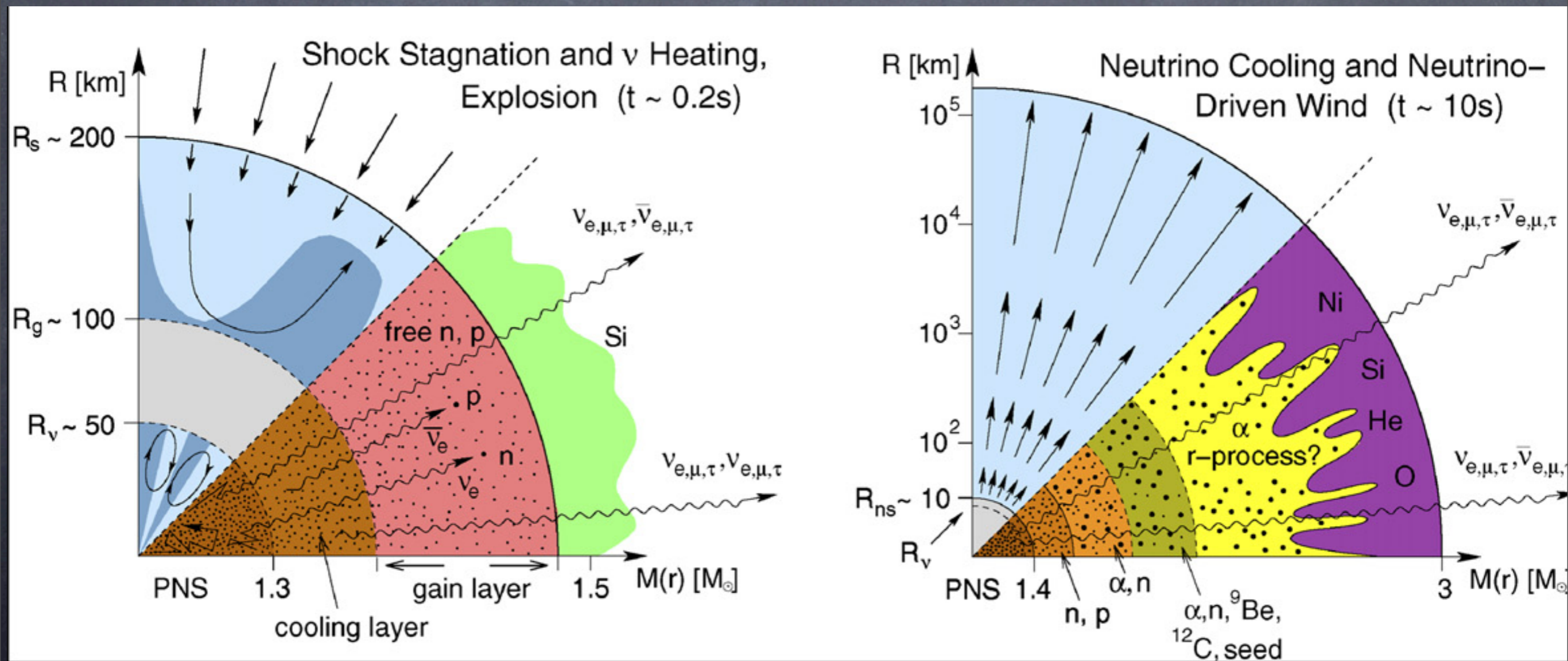
Core Collapse Basics

$$\rho_0 \sim 10^{14} \text{ g cm}^{-3}$$



Janka (2007)

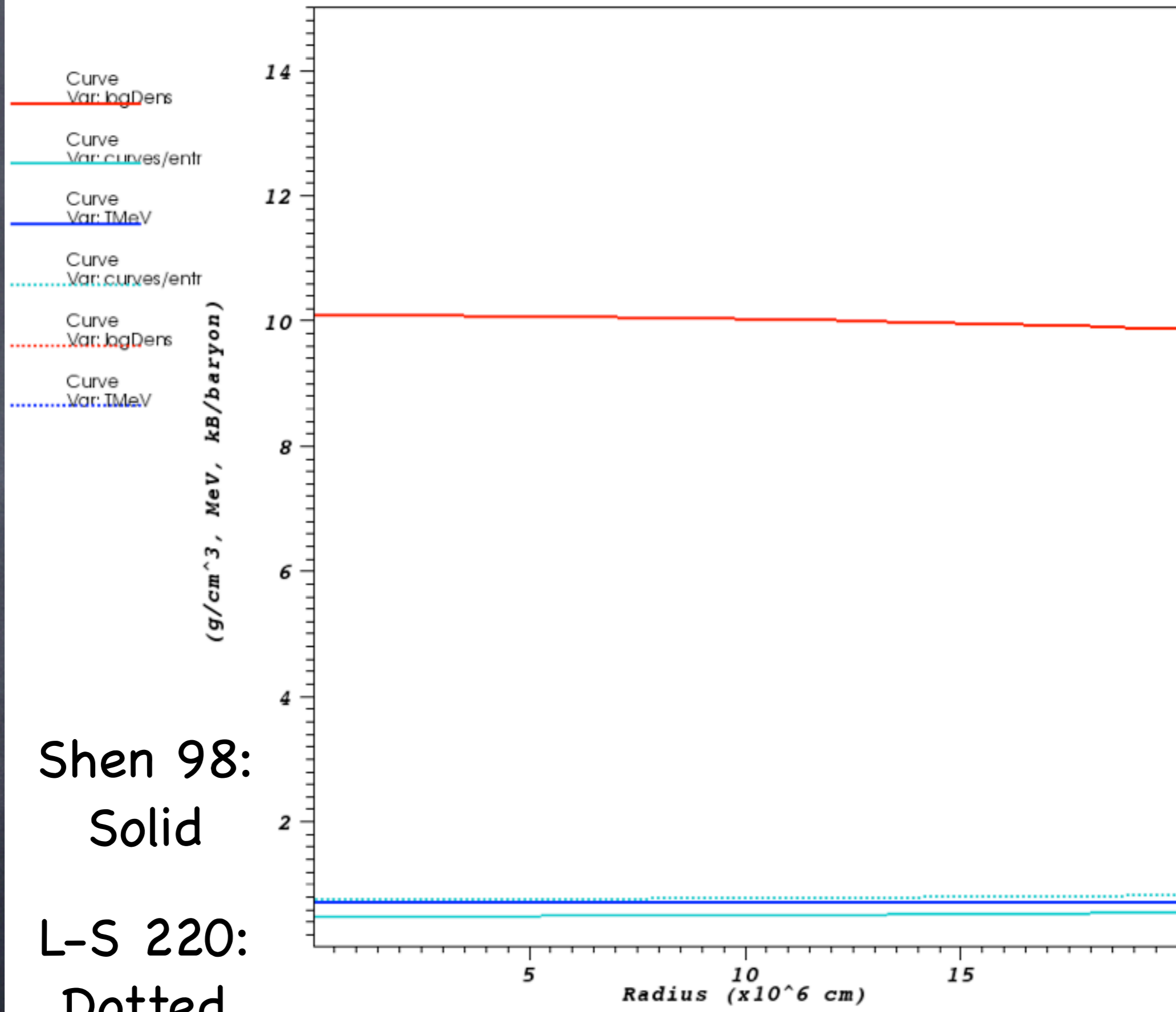
Core Collapse Basics



Janka (2007)

Shen 98:
Solid

L-S 220:
Dotted



Shen 98:
Solid

L-S 220:
Dotted



Time=0

user: smc
Mon Jul 25 15:14:29 2011

What's the big deal?

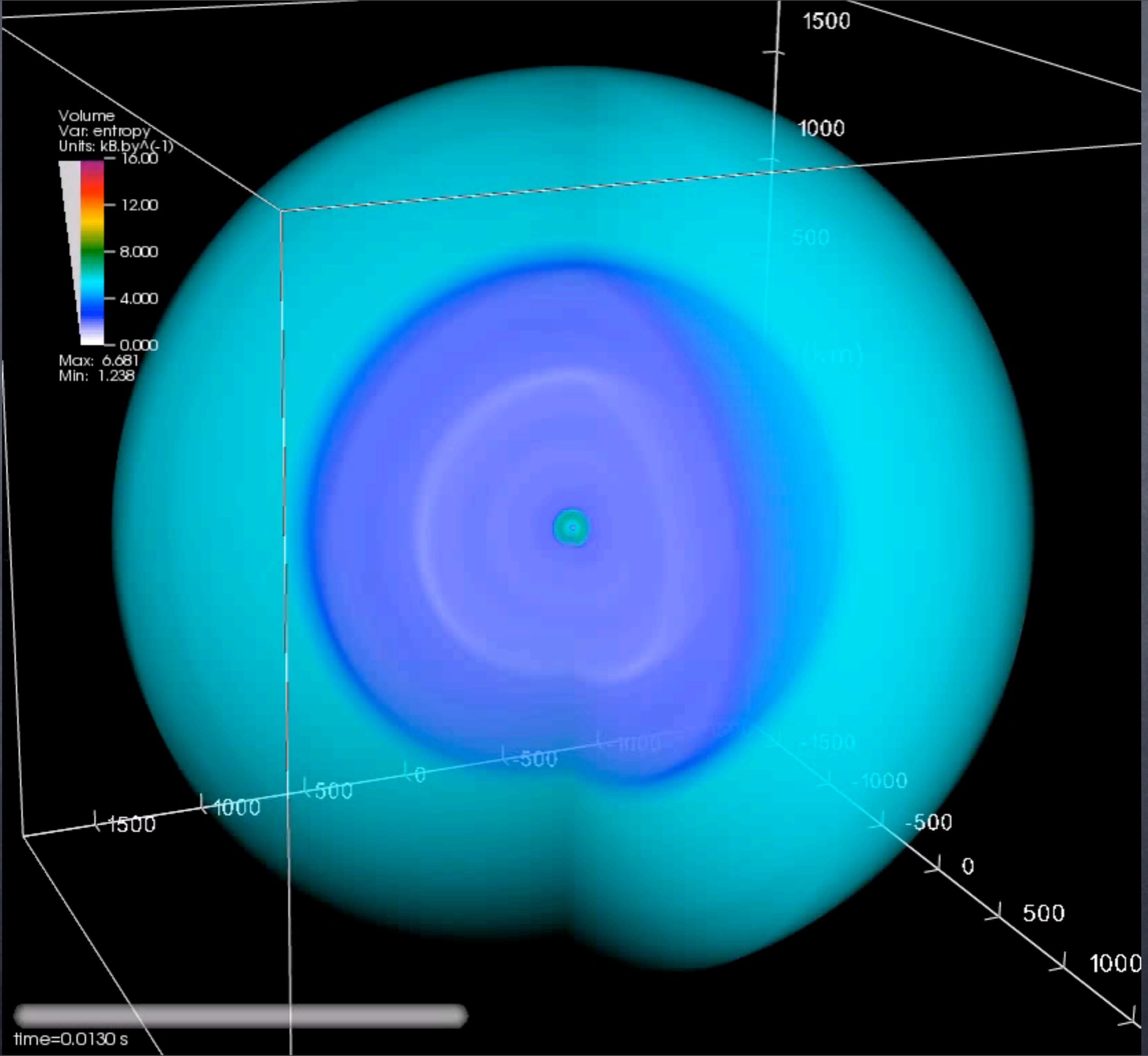
- Neutrino radiation hydrodynamics is hard!
- Highly non-local problem.
- Must have closure scheme for radiation transport (or Boltzmann) moment equations.
- Common closure schemes: Flux-limited diffusion, Ray-by-ray spectral transport, full multi-angle spectral transport, etc.

An Explosion in 2D:

- “Ray-by-ray plus” method (Buras et al. 2006)
- variable Eddington factor technique
- rather under-energetic explosion only for $11.2 M_{\text{sun}}$ star.

Buras, Marek, Janka, Rammp

Buras, Marek, Janka, Ramm



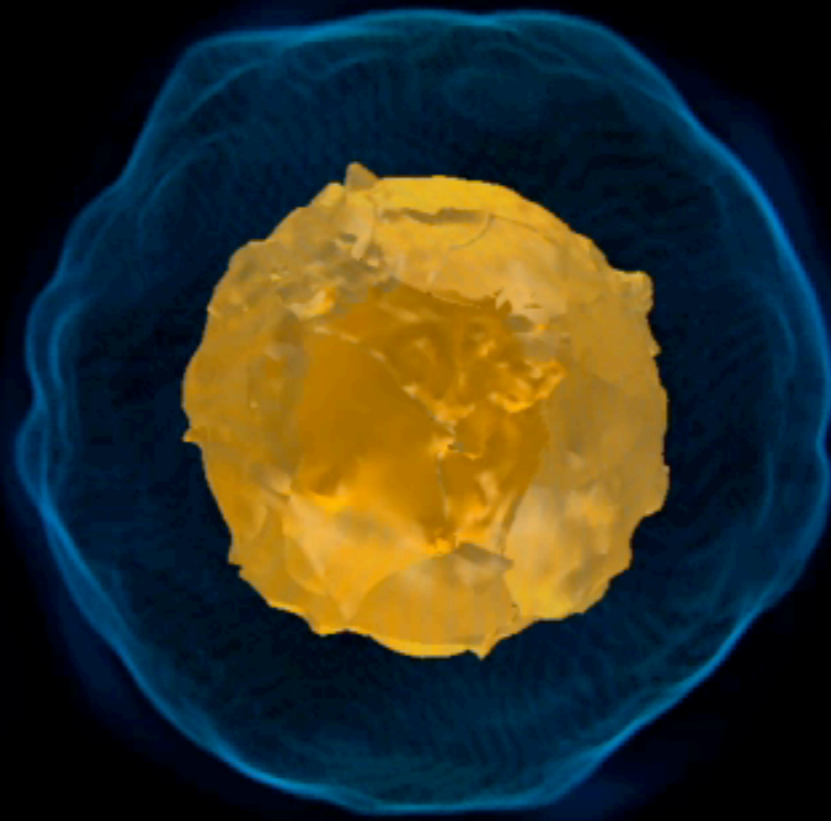
Nature isn't two dimensional...

- Murphy & Burrows (2008) and Nordhaus et al. (2010) found that explosions are more easily obtained in higher dimension.
- Parameterized neutrino heating & cooling with approximate deleptonization scheme.

Entropy

Nordhaus, Burrows, Bell, Almgren, Chupa

Entropy



Nordhaus, Burrows, Bell, Almgren, Chupa

Some Salient Features of the Neutrino Mechanism

- 10^{53} erg released in neutrinos!
- Sophisticated 1D simulations don't give explosions for progenitors bigger than about $10 M_{\text{sun}}$.
- 2D simulations only give explosions for $11.2 M_{\text{sun}}$ progenitors.
- Multidimensional effects (convection, SASI) critical to success.
- Easier to get explosion in 3D!

Neutrino Mechanism

Report Card

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Magnetorotational SNe

- All stars rotate and have magnetic fields!
- Magnetic fields can tap the energy in differential rotation to power outflows
- Some progenitors may rotate fast enough to power a magnetorotational explosion
- See Wheeler et al. 2000, 2002

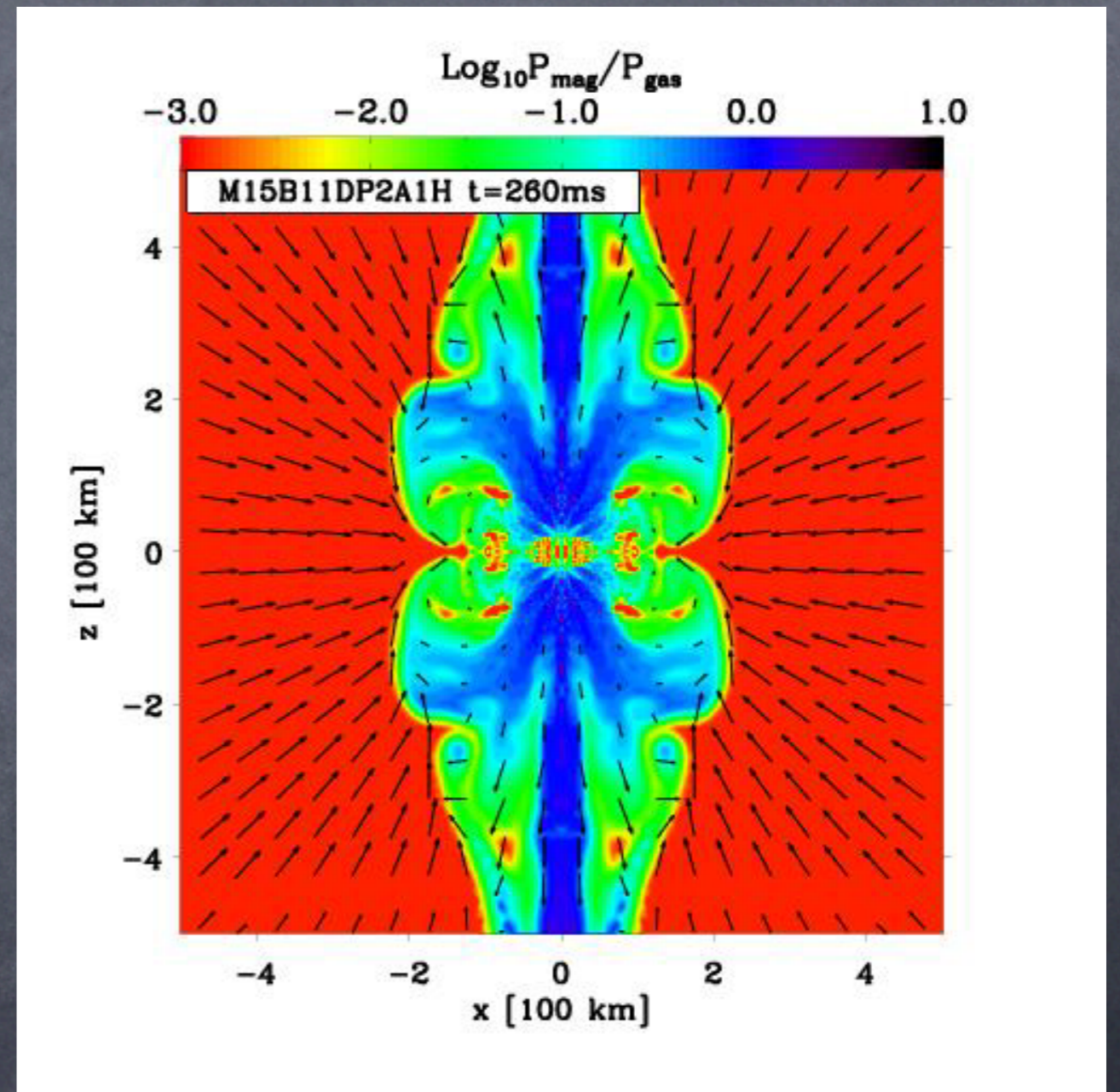
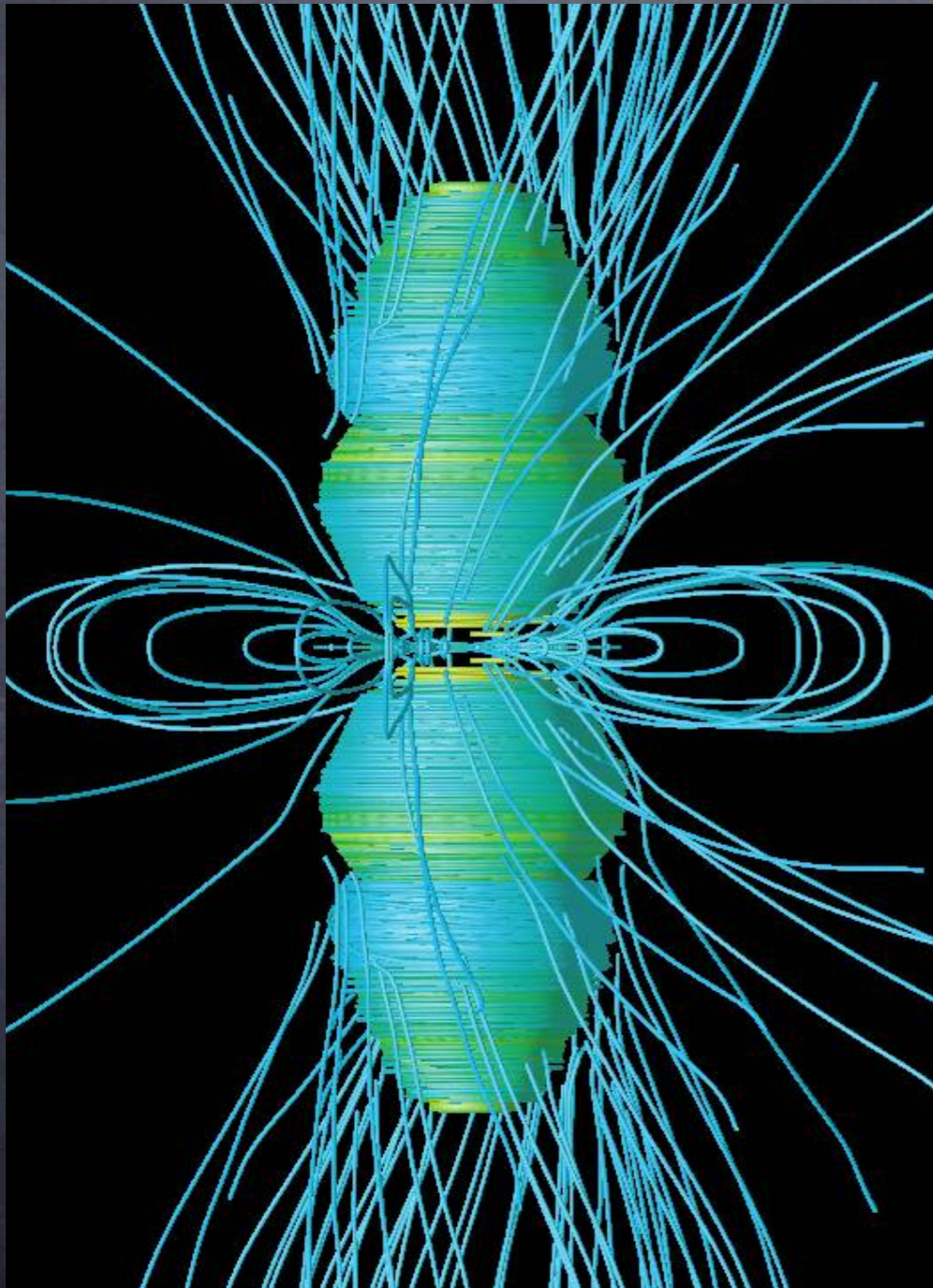
$$E_{\text{rot,PNS}} \approx \frac{1}{2} I_{\text{PNS}} \Omega_{\text{PNS}}^2$$
$$\approx 9 \times 10^{50} \text{ ergs} \left(\frac{M_{\text{PNS}}}{1.5 M_{\odot}} \right) \left(\frac{\Omega_{\text{PNS}}}{250 \text{ s}^{-1}} \right)^2 \left(\frac{R_{\text{PNS}}}{50 \text{ km}} \right)^2$$

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Rapid rotation, but in line with PNS rotation speeds
from Burrows et al. (2007).

Not all of this energy will be available to drive an
explosion!

Magnetorotational SNe



Burrows et al. 2007

Magnetorotational SNe

- Rapid rotation required for MHD jet-driven explosion
- MHD & rotation may still be important in slower rotators
- SN progenitor core rotation and B-field not well defined
- MRI not resolved in simulations

Magnetorotational SNe

- Elongated explosions
- Non-axisymmetries via instabilities
- High-velocity nickel clumps
- Complex, large-scale structures
- Mechanism for pulsar kicks
- Continuum to GRBs with higher rotation

Can Bipolar Explosions
Explain Observations?

Jet-driven Type IIP SNe

SMC, Wheeler, Milosavljević 2009, ApJ, 696, 953

- 15 M_{\odot} Red Supergiant progenitor star
- FLASH hydrodynamics code
- Jets introduced at inner boundary
- Four physically-motivated jet models
- Evolved to 500,000 seconds

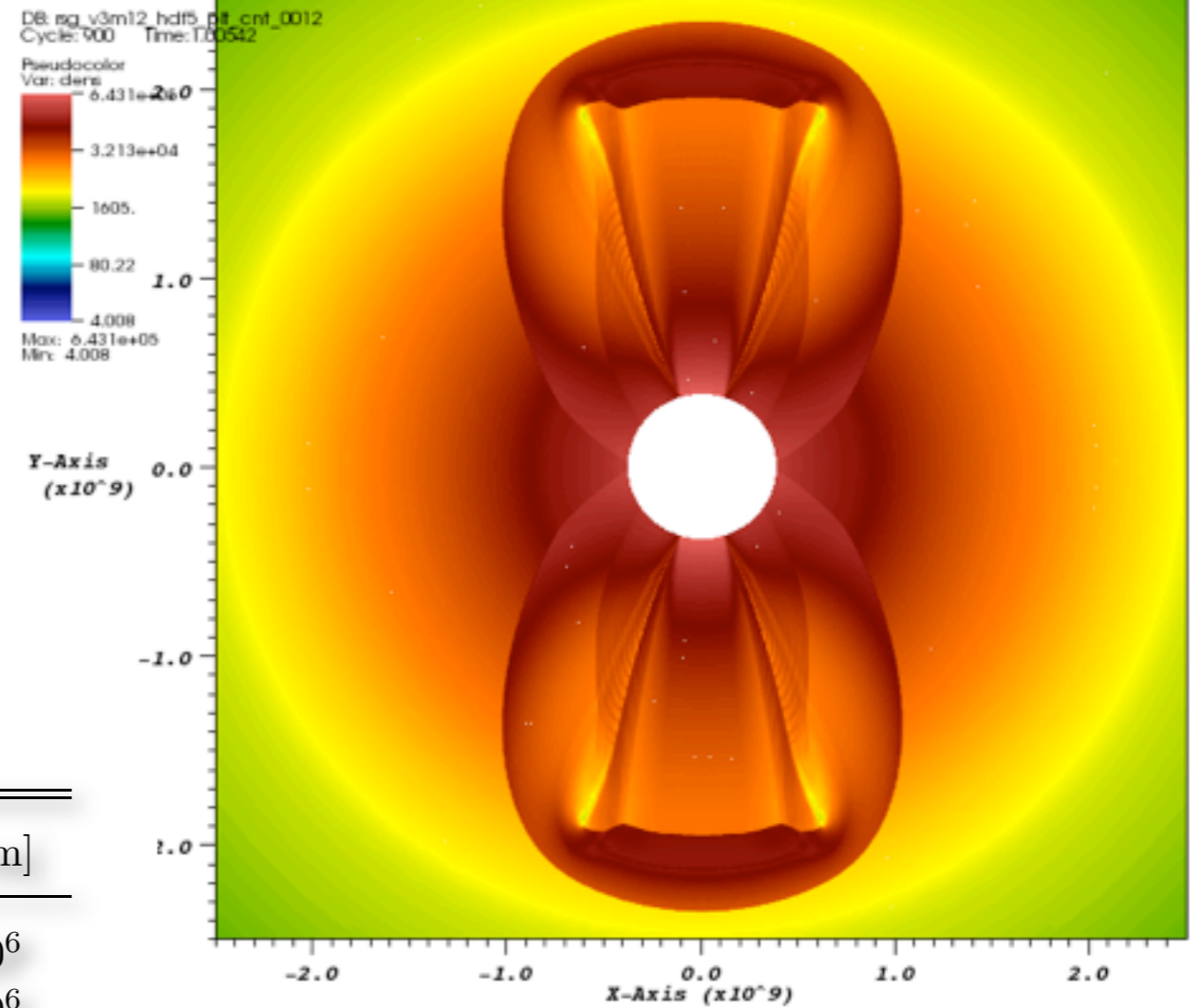
Dynamic Range

- $\Delta_r \sim 10^6$ cm out to 10^{15} cm
- 2D spherical geom.
- 7 refinement levels, re-gridding algorithm

| Stage | t_i [s] | t_f [s] | r_{in} [cm] | r_{out} [cm] | $N_{r,0}$ | Δr_{min} [cm] |
|-------|-----------------|-----------------|----------------------|----------------------|-----------|-----------------------|
| 0: | 0 | 5 | 3.82×10^8 | 3.2×10^{10} | 192 | 2.6×10^6 |
| 1: | 5 | 25 | 3.82×10^8 | 7.5×10^{10} | 192 | 6.1×10^6 |
| 2: | 25 | 100 | 1.0×10^9 | 2.5×10^{11} | 192 | 2.0×10^7 |
| 3: | 100 | 300 | 2.0×10^9 | 9.0×10^{11} | 192 | 7.3×10^7 |
| 4: | 300 | 1×10^3 | 4.0×10^9 | 3.0×10^{12} | 320 | 1.5×10^8 |
| 5: | 1×10^3 | 3×10^3 | 1.0×10^{10} | 6.0×10^{12} | 320 | 2.9×10^8 |
| 6: | 3×10^3 | 1×10^4 | 2.0×10^{10} | 2.0×10^{13} | 352 | 8.9×10^8 |
| 7: | 1×10^4 | 3×10^4 | 5.0×10^{10} | 6.0×10^{13} | 352 | 2.7×10^9 |
| 8: | 3×10^4 | 1×10^5 | 1.0×10^{11} | 2.0×10^{14} | 384 | 8.1×10^9 |
| 9: | 1×10^5 | 2×10^5 | 2.0×10^{11} | 4.0×10^{14} | 384 | 1.6×10^{10} |
| 10: | 2×10^5 | 5×10^5 | 5.0×10^{11} | 1.0×10^{15} | 384 | 4.1×10^{10} |

Dynamic Range

- $\Delta_r \sim 10^6$ cm out to 10^{15} cm
- 2D spherical geom.
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user: smc
Tue Apr 15 12:04:28 2008

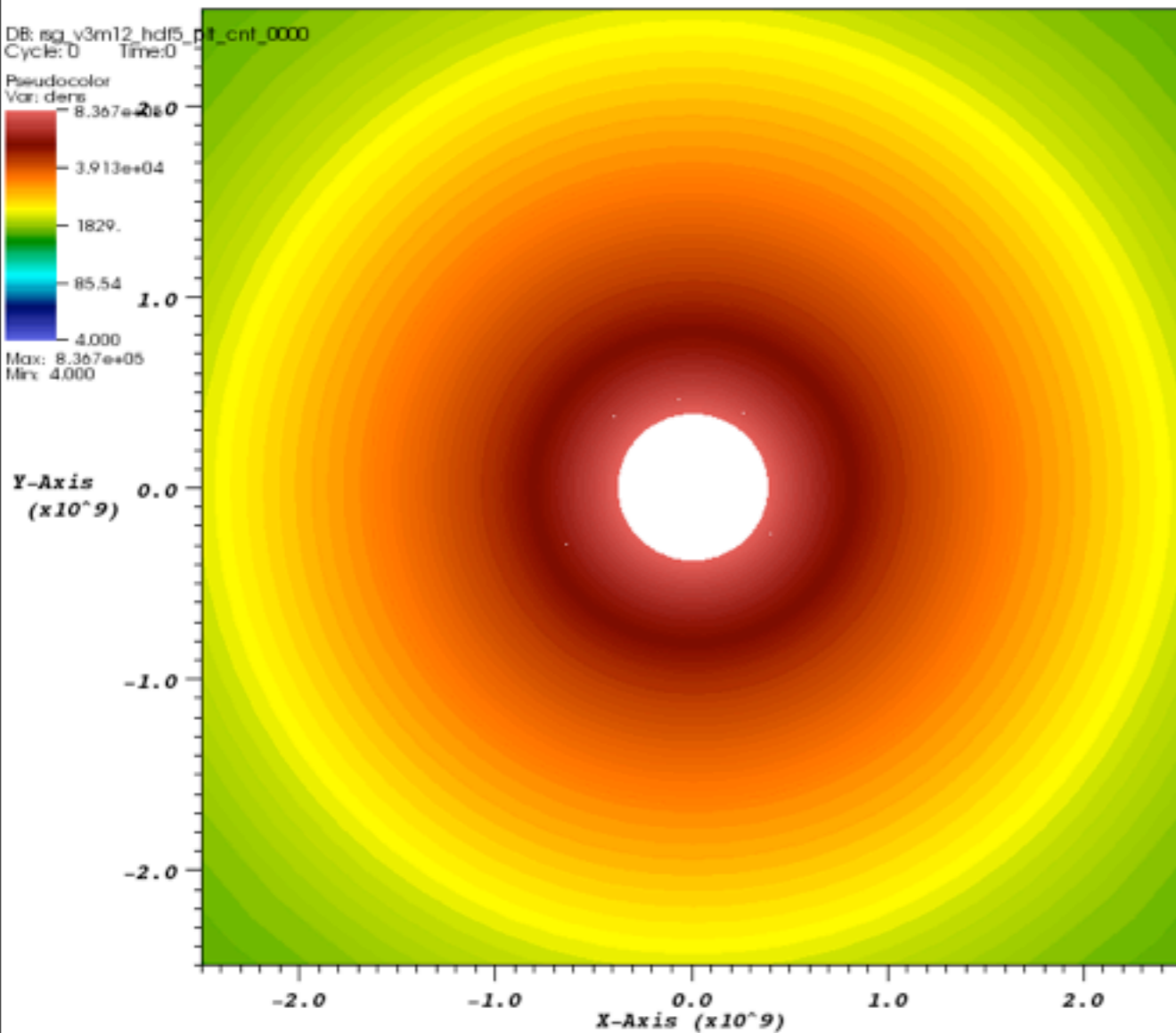
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Density

Fast, kinetic

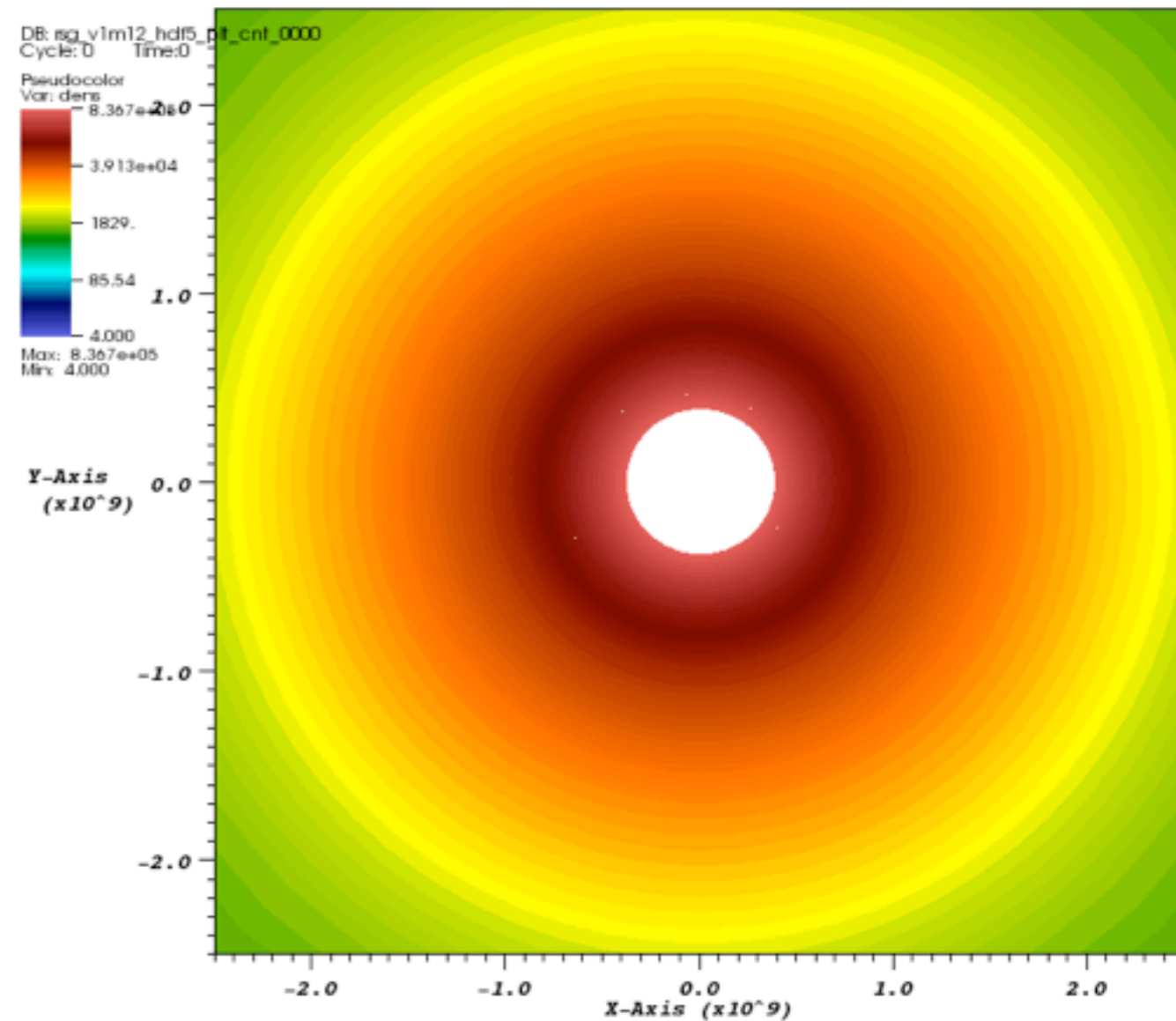
Slow, thermal

Density



user: smc
Tue Apr 1 11:56:12 2008

Fast, kinetic



user: smc
Wed Apr 9 10:46:35 2008

Slow, thermal

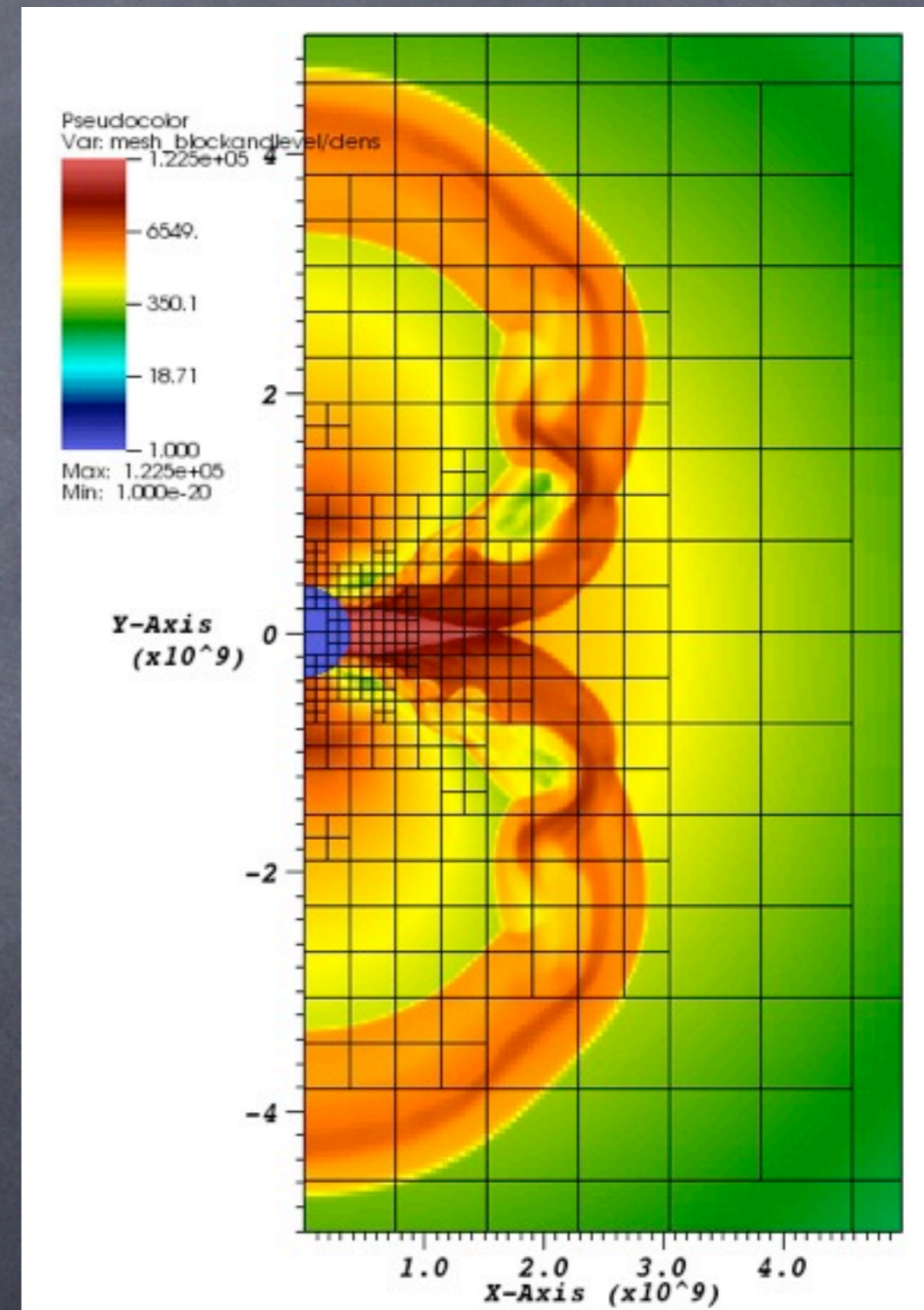
Jet-driven Type Ib SNe

SMC, Pooley, Wheeler, Milosavljević 2011, ApJ, 727, 104

- 2.5 & 6 M_{\odot} helium core progenitors
- Thermal & kinetic jet models in each
- FLASH hydro
- Custom post-processing radiation modeling

Dynamic range

- 2D cylindrical geom.
- Radius, time-dependent max. refinement level
- Modified FLASH to excise central hole
- Hole radius expands with time
- No need for regrid; start with 25 refinement levels



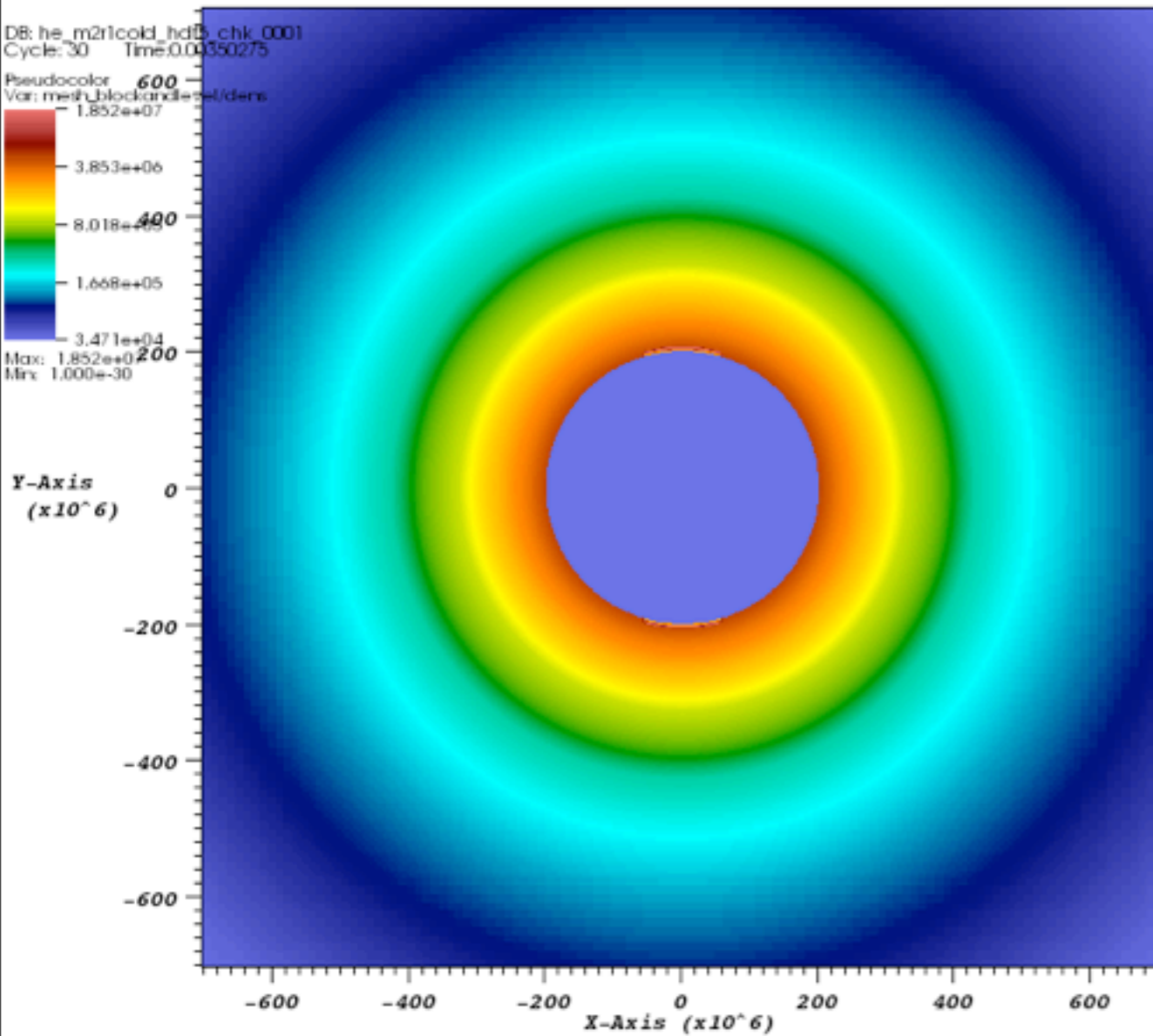
Density

Kinetic

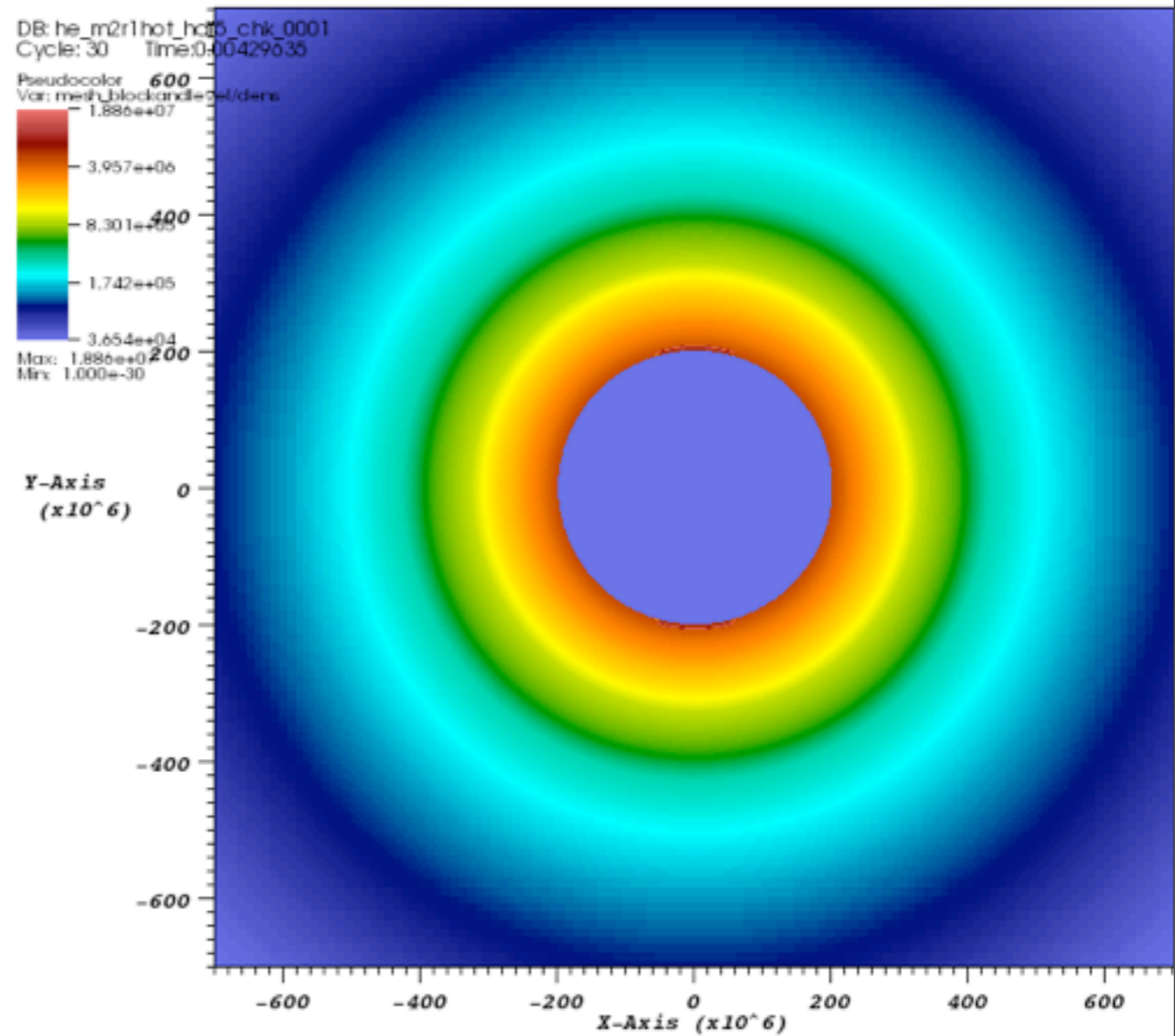
Thermal

Smaller Progenitor

Density



User: smc
Mon Sep 28 11:56:02 2011



User: smc
Tue Sep 29 17:09:10 2011

Kinetic

Thermal

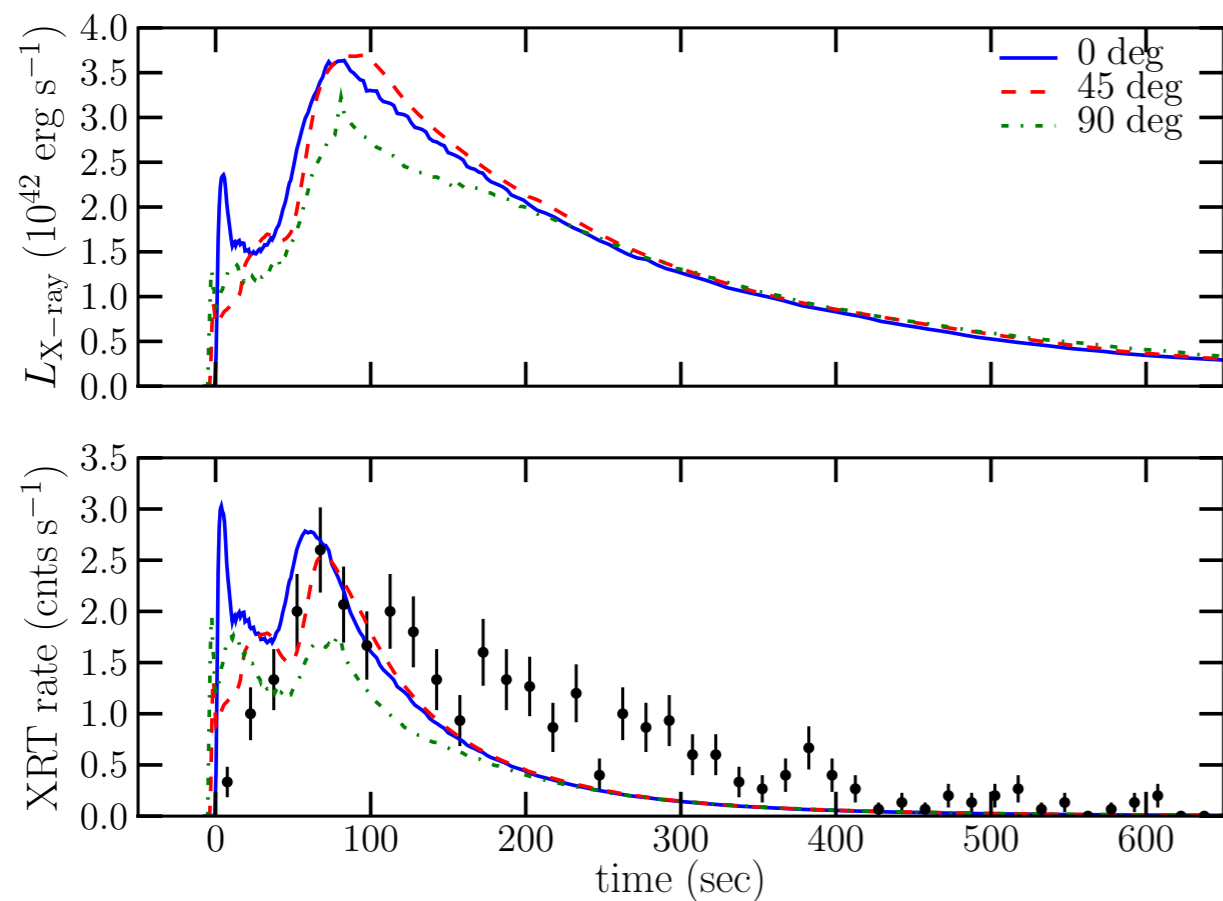
Smaller Progenitor

Small progenitor, kinetic jets

- LC shape and time scales about right

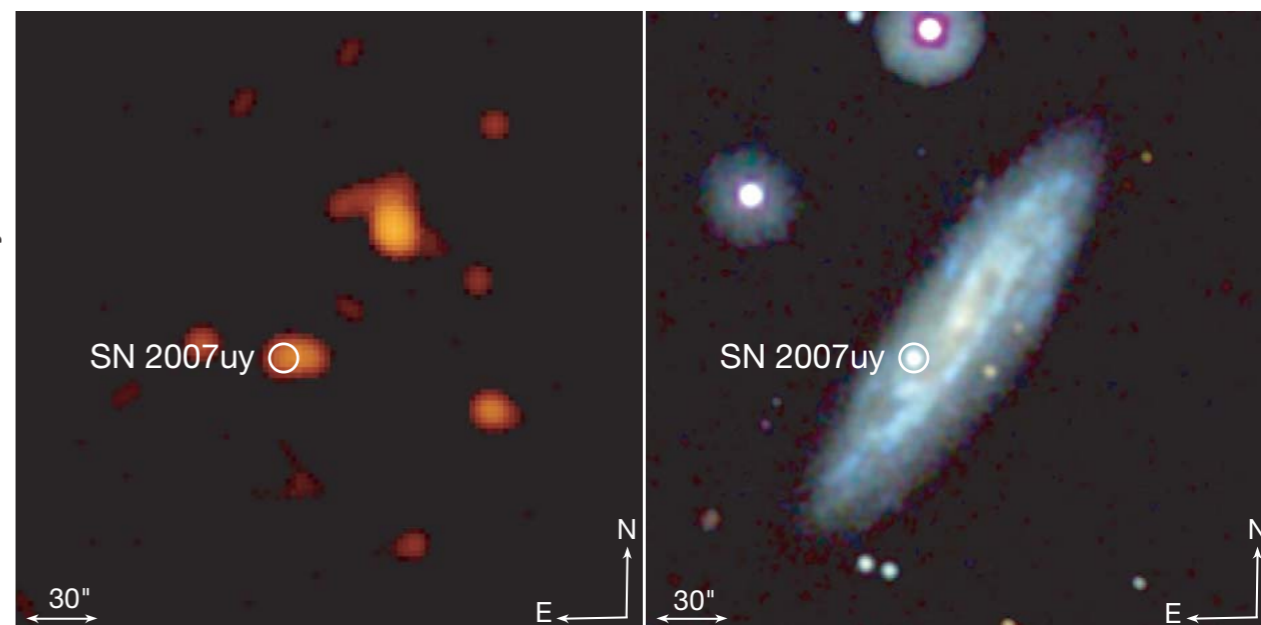
X-ray

UV



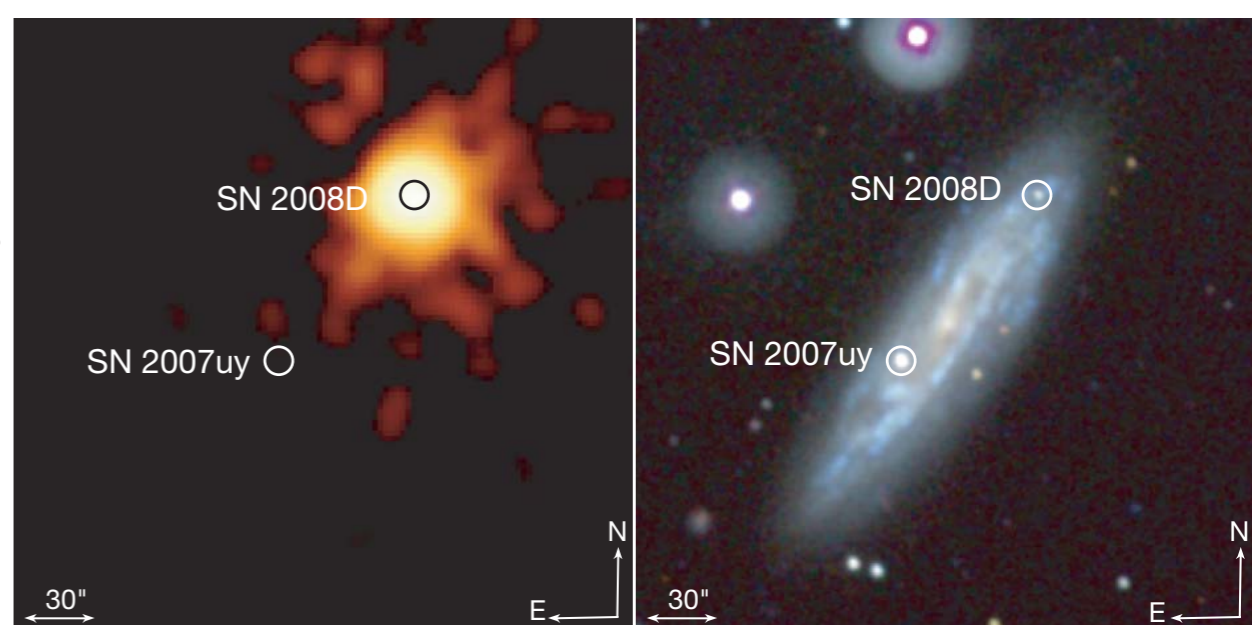
a

2008 January 7



b

2008 January 9



Soderberg et al. 2008

- Spectra are too soft

- LC time scales right-on!

- Phenomenological, jet-driven explosions can explain many of the observations, but lack crucial physics.
- Need nuclear EOS, neutrinos, and MHD:

Euler Equations

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Mass

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) + \nabla P = \rho \mathbf{g}$$

Momentum

$$\frac{\partial \rho E}{\partial t} + \nabla \cdot [(\rho E + P) \mathbf{v}] = \rho \mathbf{v} \cdot \mathbf{g}$$

Energy

$$E = \epsilon + \frac{1}{2} |\mathbf{v}|^2$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

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$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B}) + \nabla P_* = \rho \mathbf{g} + \nabla \cdot \boldsymbol{\tau}$$

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$$P_* = P + \frac{B^2}{2}$$

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$$P_* = P + \frac{B^2}{2} \quad \boldsymbol{\tau} = \text{viscosity tensor}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

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$\boldsymbol{\tau}$ = viscosity tensor

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$$\begin{aligned} \frac{\partial \rho E}{\partial t} + \nabla \cdot [(\rho E_* + P_*) \mathbf{v} - \mathbf{B}(\mathbf{v} \cdot \mathbf{B})] &= \rho \mathbf{v} \cdot \mathbf{g} + \nabla \cdot (\mathbf{v} \cdot \boldsymbol{\tau} + \sigma \nabla T) \\ &+ \nabla \cdot [\mathbf{B} \times (\eta \nabla \times \mathbf{B})] \end{aligned}$$

$\boldsymbol{\tau}$ = viscosity tensor

$$P_* = P + \frac{B^2}{2}$$

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$$P_* = P + \frac{B^2}{2}$$

$$E_* = \epsilon + \frac{1}{2} v^2 + \frac{1}{2} \frac{B^2}{\rho}$$

$\boldsymbol{\tau}$ = viscosity tensor
 σ = thermal conductivity
 η = resistivity

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B}) + \nabla P_* = \rho \mathbf{g} + \nabla \cdot \boldsymbol{\tau}$$

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Induction Equation:

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v}) = -\nabla \times (\eta \nabla \times \mathbf{B})$$

| | | |
|---|---------------------|------------------------|
| $P_* = P + \frac{B^2}{2}$ | $\boldsymbol{\tau}$ | = viscosity tensor |
| $E_* = \epsilon + \frac{1}{2} v^2 + \frac{1}{2} \frac{B^2}{\rho}$ | σ | = thermal conductivity |
| | η | = resistivity |

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B}) + \nabla P_* = \rho \mathbf{g} + \nabla \cdot \tau$$

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| | | |
|---|----------|------------------------|
| $P_* = P + \frac{B^2}{2}$ | τ | = viscosity tensor |
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$$\nabla \cdot \mathbf{B} = 0$$

Solenoidal Constraint

Three Common Methods to Satisfy Constraint:

1. Elliptic projection (Brackbill & Barnes 1980)
2. Constrained transport (Evans & Hawley 1988)
3. Divergence cleansing (Powell et al. 1999)

B-field Amplification in CCSNe

- Field compression: field carried along with collapsing plasma: “flux-freezing”
- Field winding: linear process, wraps up field lines. $B_\phi \approx 2\pi n_\phi B_p$
- Magnetorotational Instability (MRI): exponential growth of initial field. Saturation field strengths as high as 10^{15} – 10^{16} G.

MRI

$$\tau_{\text{MRI}} \sim 4\pi \left(\frac{\partial \ln r}{\partial \Omega} \right) \sim 2P$$

$$\lambda_{\text{MRI}}^{\text{max}} \sim \frac{2\pi v_A}{\Omega} \sim (10^4 \text{ cm}) P_{10} \frac{B_{12}}{\rho_{11}^{1/2}}$$

- Grows on the rotational time scale.
- Requires restrictive resolution! 100 times that of high-resolution CCSNe sims.
- See, e.g., Akiyama et al. (2003), Obergaulinger et al. (2011).

Conclusions

- Observations show that CCSN are aspherical
- Physics of CCSN & the Neutrino Mechanism
- Robust neutrino-driven explosions are not found in sophisticated calculations
- Rotation and magnetic fields may play an important role in shaping or driving CCSN explosions
- Bipolar explosions may explain observations indicating asymmetry
- MHD effects are important in CCSN

Extra Slides

Basics of Unsplit Staggered-Mesh (USM) Constrained Transport

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x} + \frac{\partial \mathbf{G}}{\partial y} = 0$$

$$\mathbf{U} = (\rho, \rho u, \rho v, \rho w, B_x, B_y, B_z, E)^T$$

$$\mathbf{F} = \begin{pmatrix} \rho u \\ \rho u^2 + p_{\text{tot}} - B_x^2 \\ \rho uv - B_y B_x \\ \rho uw - B_z B_x \\ 0 \\ uB_y - vB_x (= -E_z) \\ uB_z - wB_x (= E_y) \\ (E + p_{\text{tot}})u - B_x(uB_x + vB_y + wB_z) \end{pmatrix}, \quad \mathbf{G} = \begin{pmatrix} \rho v \\ \rho vu - B_x B_y \\ \rho v^2 + p_{\text{tot}} - B_y^2 \\ \rho vw - B_z B_y \\ vB_x - uB_y (= E_z) \\ 0 \\ vB_z - wB_y (= -E_x) \\ (E + p_{\text{tot}})v - B_y(uB_x + vB_y + wB_z) \end{pmatrix}.$$

From Lee & Deane (2009)

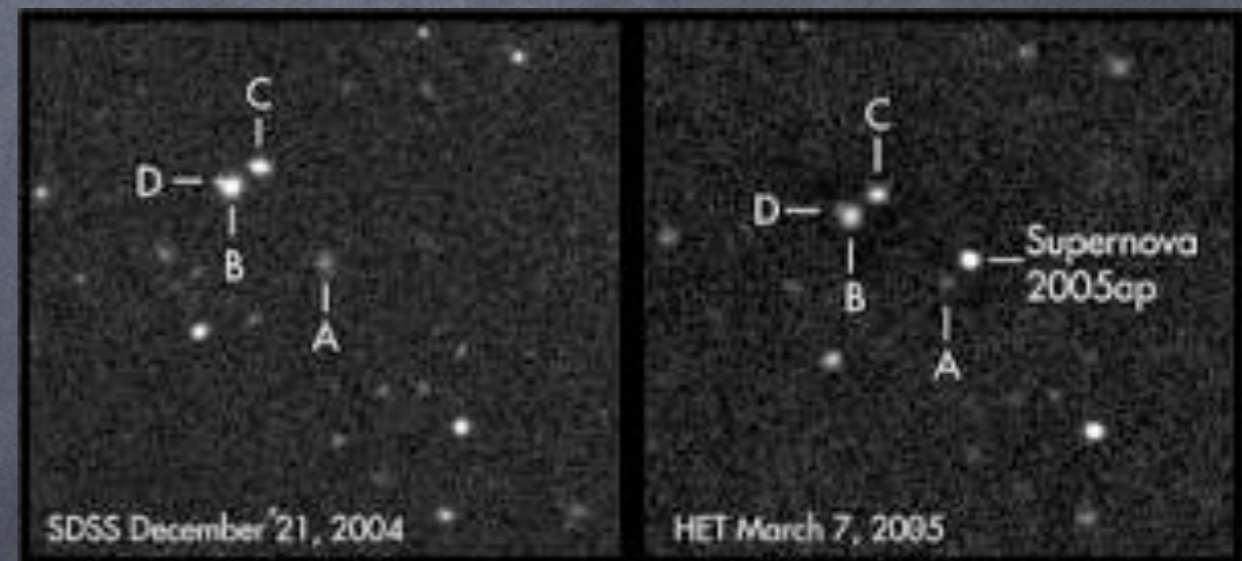
MHD Waves

- Alfvén waves

-

SN Polarization

- Cannot “see” the shapes of distant SNe
- Can get wavelength-dependent info on the shapes of the photosphere and line-forming regions
- Measure Stokes parameters:



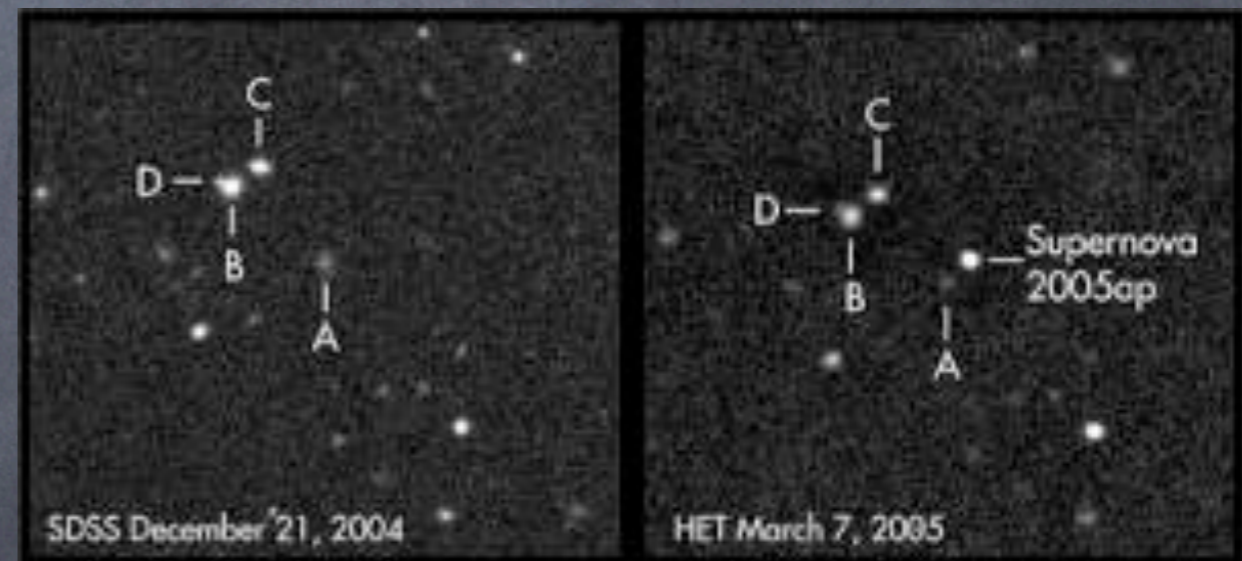
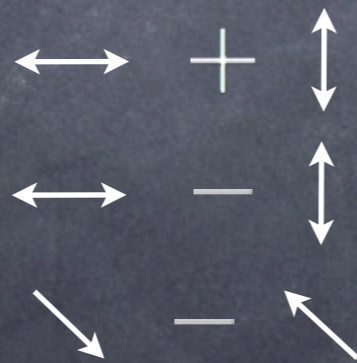
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$$I = I_0 + I_{90}$$

$$Q = I_0 - I_{90}$$

$$U = I_{45} - I_{-45}$$



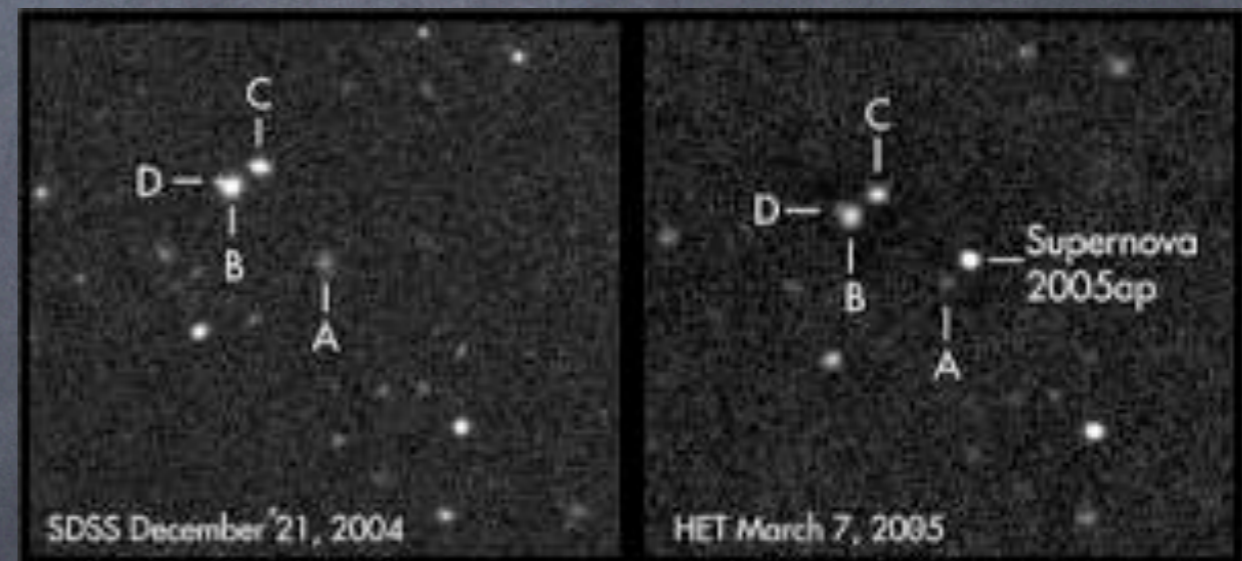
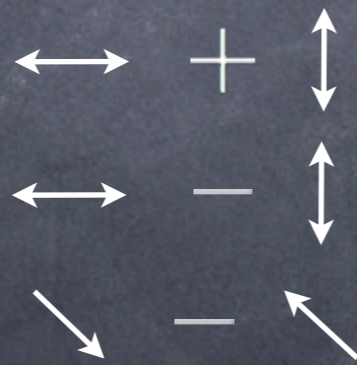
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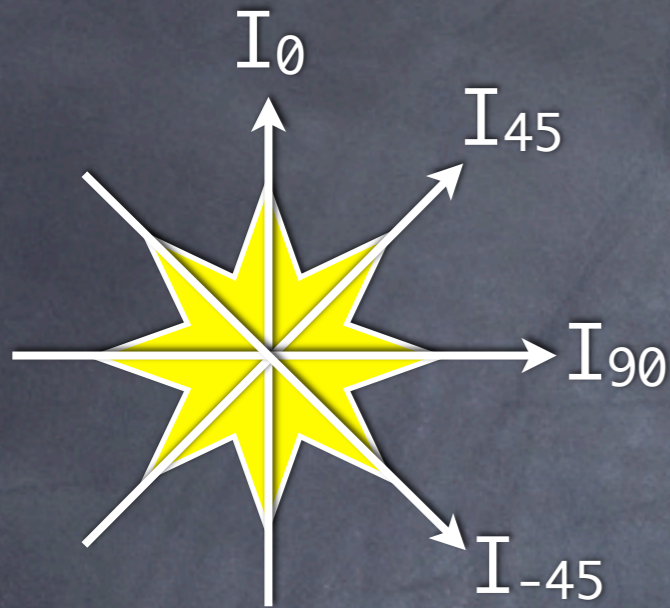


$$P = \sqrt{Q^2/I^2 + U^2/I^2} = \sqrt{q^2 + u^2} \quad \chi = \frac{1}{2} \tan^{-1}(u/q)$$

SN Polarization

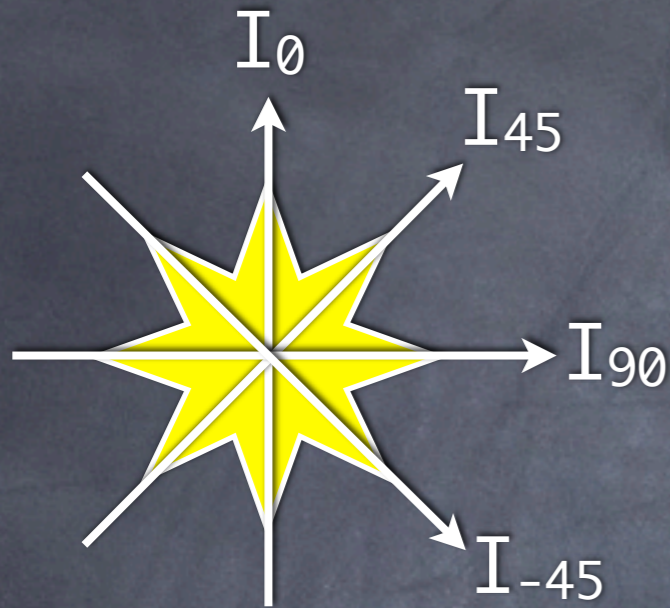


SN Polarization



- $P=Q=U=0$: no net polarization, circularly symmetric

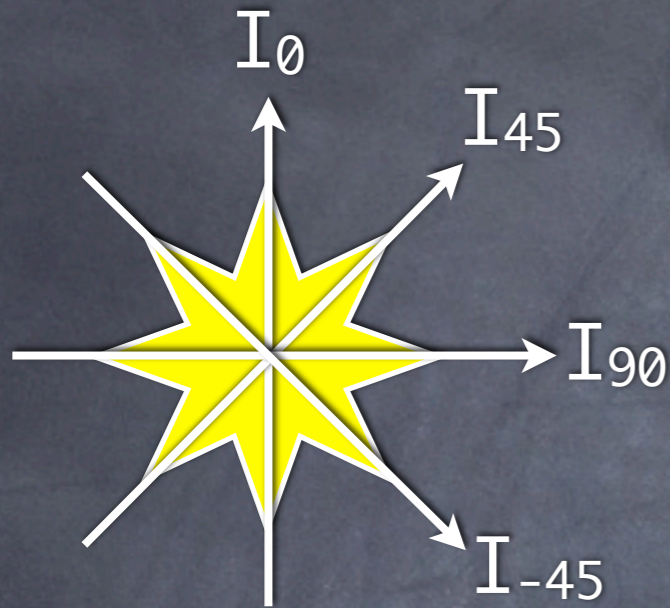
SN Polarization



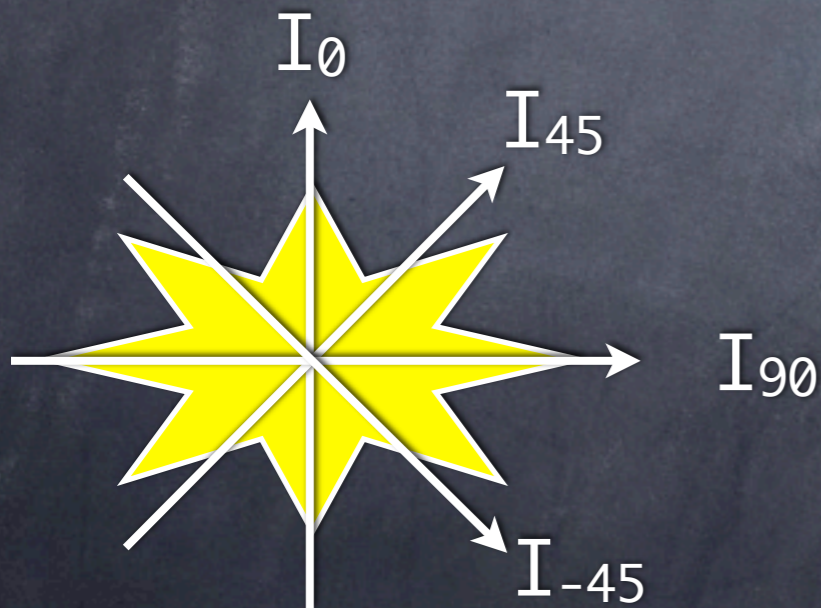
- $P=Q=U=0$: no net polarization, circularly symmetric



SN Polarization



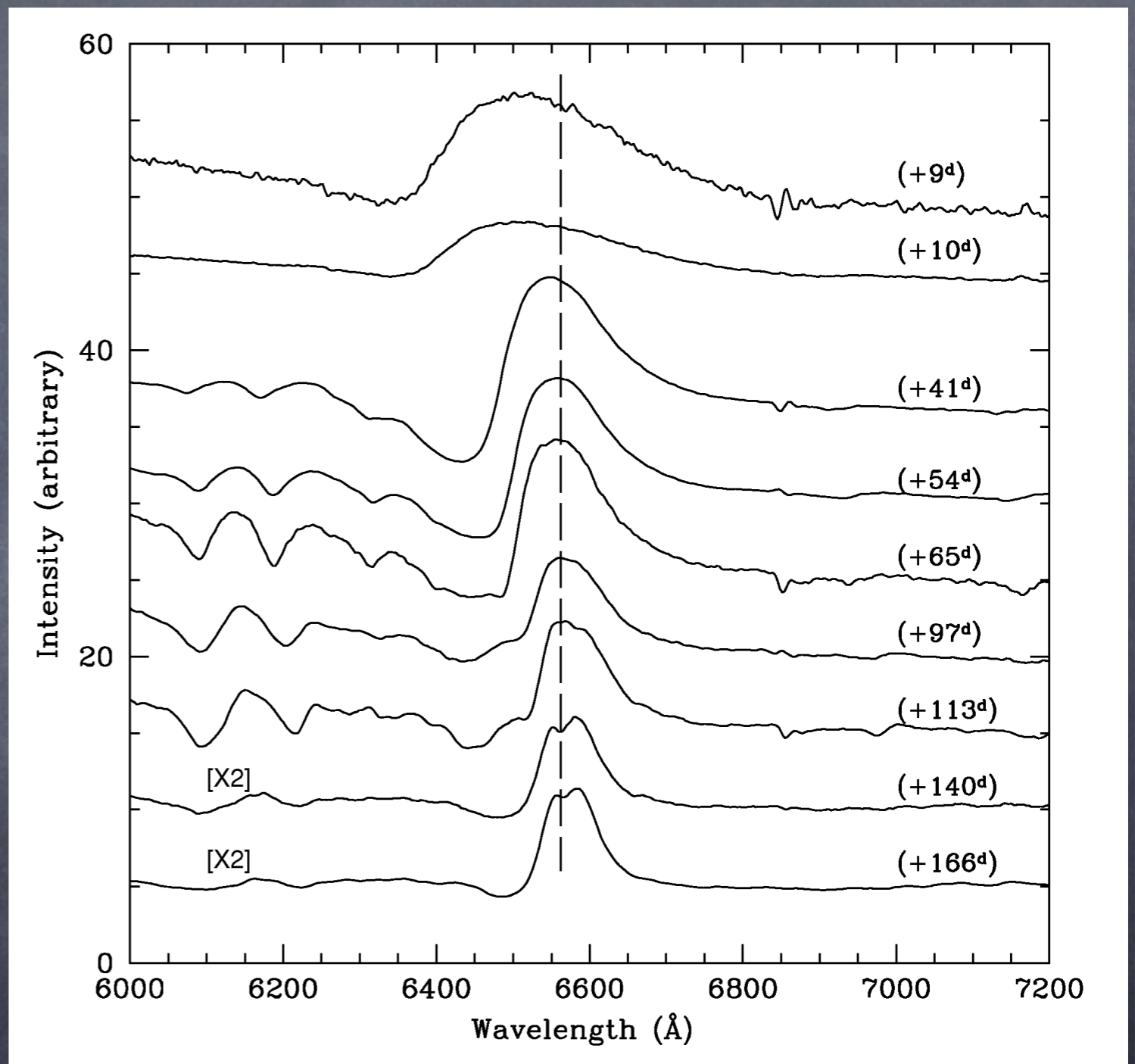
- $P=Q=U=0$: no net polarization, circularly symmetric



- $P, Q, U \neq 0$: net polarization, asymmetric emitting region

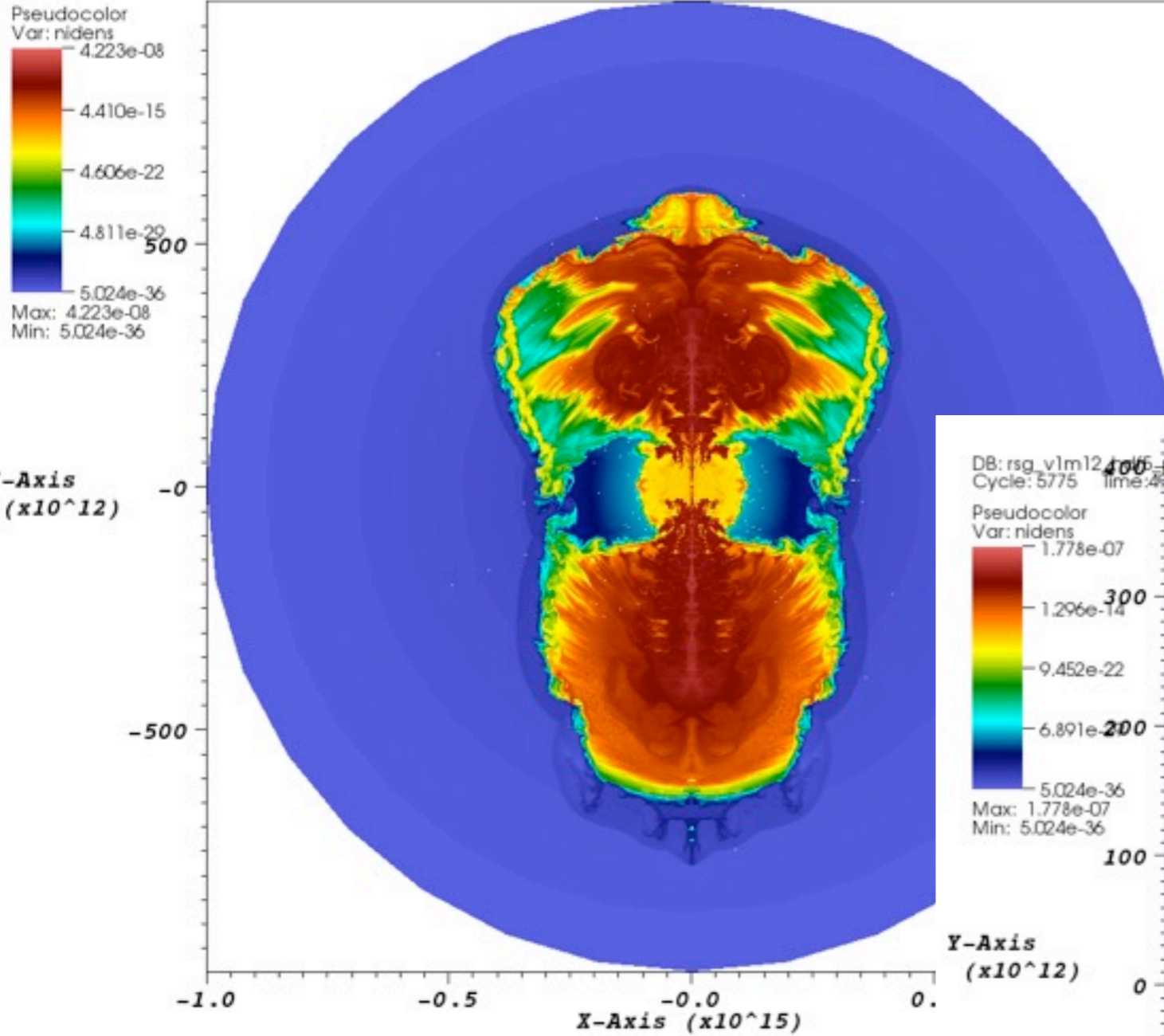
H α sub-structure

- Appearance of peaks in line require fast nickel clumps

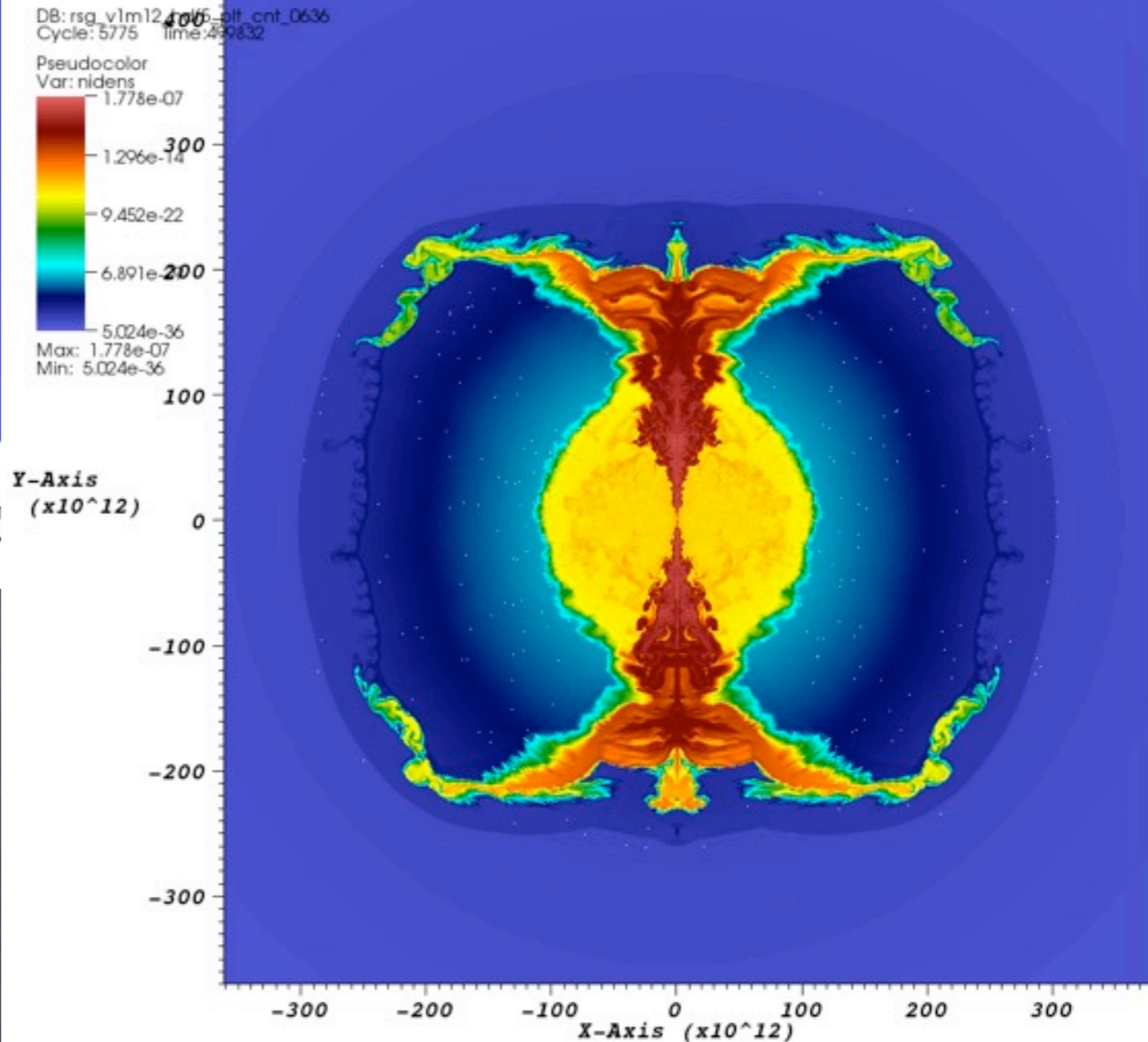


Elmhamdi et al. 2003

DB: rsg_v3m12_hdf5_plt_cnt_1123
Cycle: 10725 Time: 498793



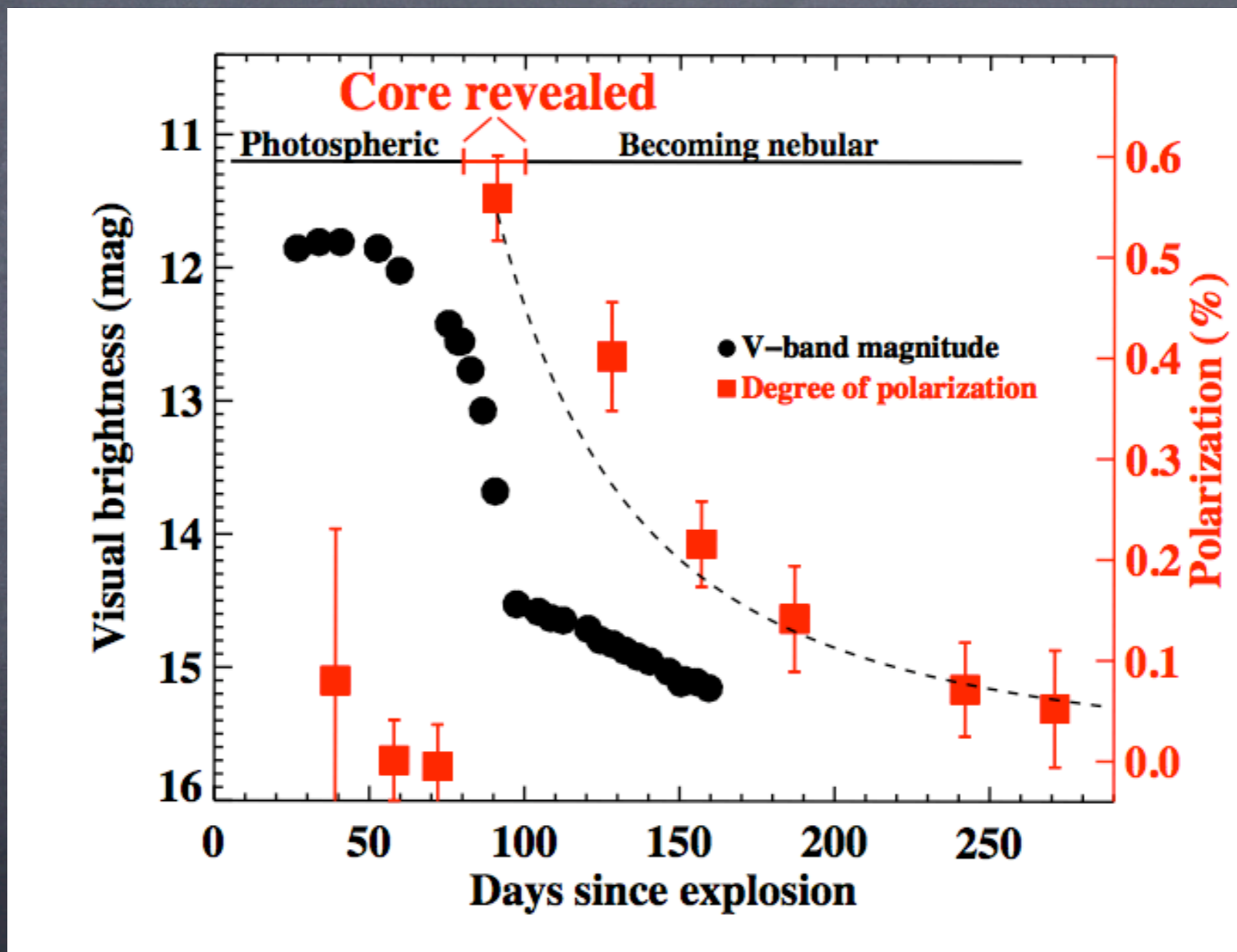
Nickel Clumps



- High-velocity bipolar clumps
- May effect spectral lines

Type IIP Polarization

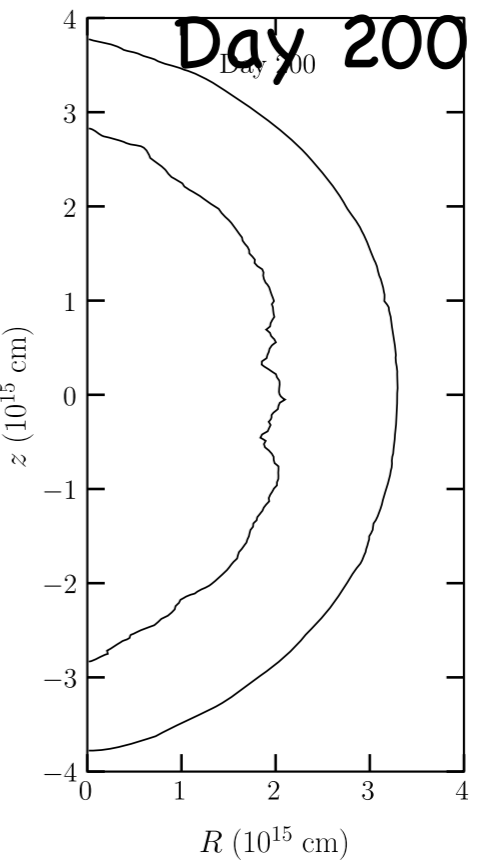
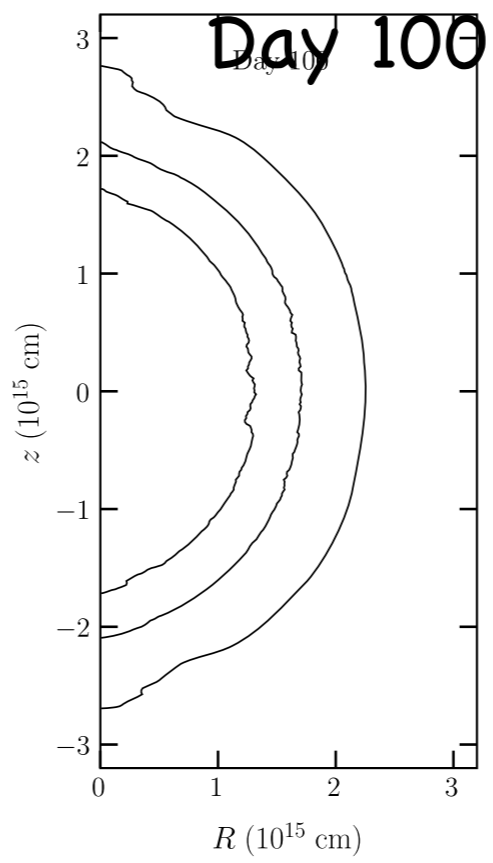
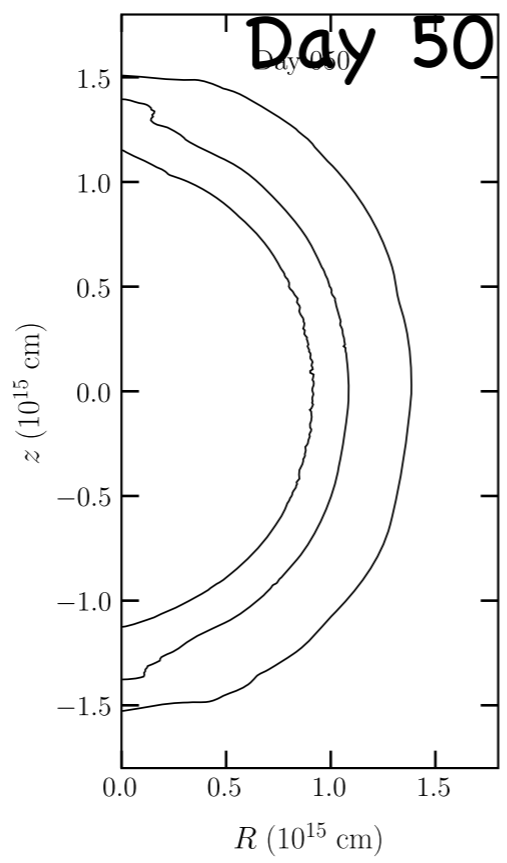
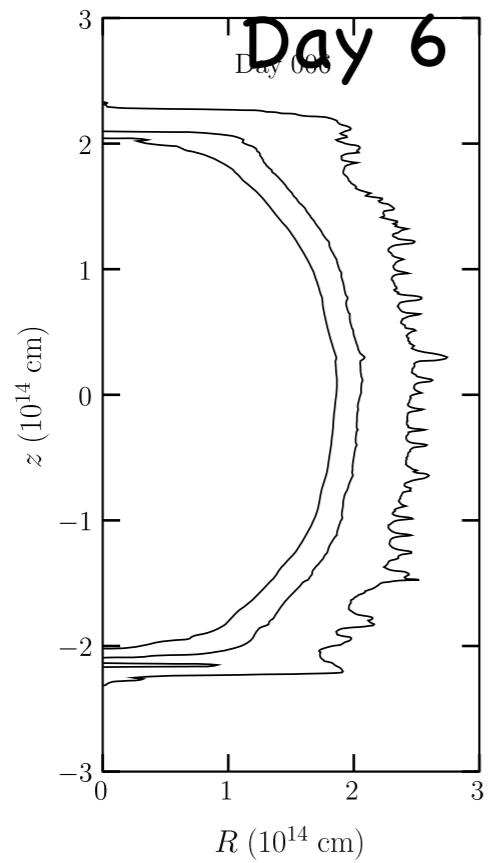
SN 2004dj



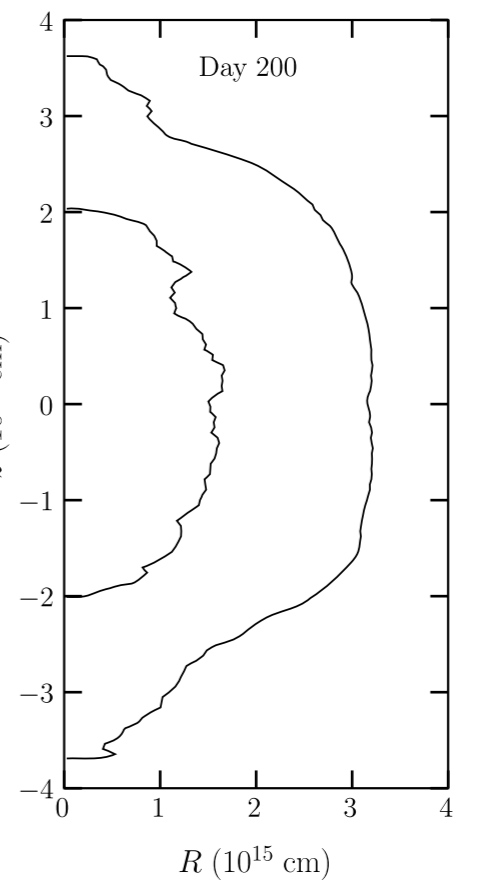
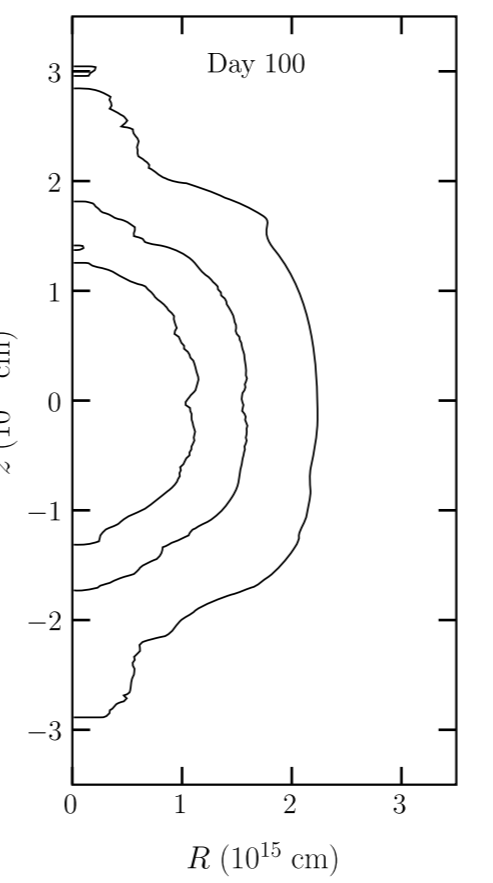
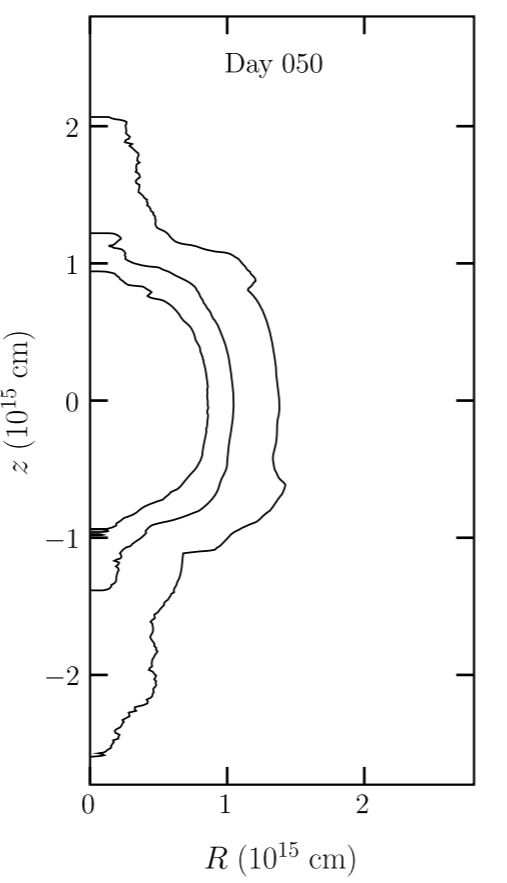
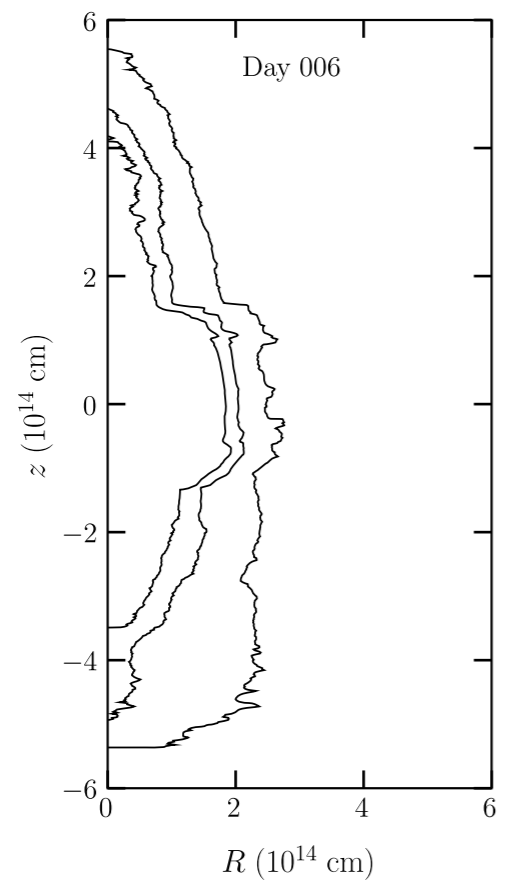
Leonard et al. 2006

Photosphere Shapes

Thermal



Kinetic



$\tau_{es} =$
1,
10,
30